

Zeitschrift: Helvetica Physica Acta
Band: 68 (1995)
Heft: 4

Artikel: Representation of a spin-1 entity as a joint system of two spin-1/2 entities on which we introduce correlations of the second kind
Autor: Coecke, Bob
DOI: <https://doi.org/10.5169/seals-116746>

Nutzungsbedingungen

Die ETH-Bibliothek ist die Anbieterin der digitalisierten Zeitschriften auf E-Periodica. Sie besitzt keine Urheberrechte an den Zeitschriften und ist nicht verantwortlich für deren Inhalte. Die Rechte liegen in der Regel bei den Herausgebern beziehungsweise den externen Rechteinhabern. Das Veröffentlichen von Bildern in Print- und Online-Publikationen sowie auf Social Media-Kanälen oder Webseiten ist nur mit vorheriger Genehmigung der Rechteinhaber erlaubt. [Mehr erfahren](#)

Conditions d'utilisation

L'ETH Library est le fournisseur des revues numérisées. Elle ne détient aucun droit d'auteur sur les revues et n'est pas responsable de leur contenu. En règle générale, les droits sont détenus par les éditeurs ou les détenteurs de droits externes. La reproduction d'images dans des publications imprimées ou en ligne ainsi que sur des canaux de médias sociaux ou des sites web n'est autorisée qu'avec l'accord préalable des détenteurs des droits. [En savoir plus](#)

Terms of use

The ETH Library is the provider of the digitised journals. It does not own any copyrights to the journals and is not responsible for their content. The rights usually lie with the publishers or the external rights holders. Publishing images in print and online publications, as well as on social media channels or websites, is only permitted with the prior consent of the rights holders. [Find out more](#)

Download PDF: 21.02.2026

ETH-Bibliothek Zürich, E-Periodica, <https://www.e-periodica.ch>

Representation of a Spin-1 Entity as a Joint System of Two Spin- $\frac{1}{2}$ Entities on Which we Introduce Correlations of the Second Kind

By Bob Coecke

TENA, Free University,
Pleinlaan 2, B-1050 Brussels, Belgium

(21.V.1995, revised 15.VII.1995)

Abstract. We present a model system in \mathbb{R}^3 for a spin-1 quantum entity. If the state of the model system corresponds to a $|s| = 1$ -state of a spin-1 quantum entity, it behaves as a joint system of two separated spin- $\frac{1}{2}$ quantum entities. If it corresponds to a $|s| = 0$ -state of a spin-1 quantum entity, it behaves as the joint system of two spin- $\frac{1}{2}$ quantum entities that are entangled by correlations of the second kind, i.e., it behaves similar as a joint system of two entangled spin- $\frac{1}{2}$ quantum entities in a singlet state.

1 Introduction

In [1] and [3], Aerts shows that the joint system of two separated quantum entities cannot be described as the projection lattice of the tensor product. As a consequence, alternative descriptions of joint systems should be reconsidered. In [4], Aerts presents a model system for a quantum system in a singlet state, subjected to Aspect-like measurements. In this model system, Aerts introduces the concept of *correlations of the second kind*, i.e., correlations created during the measurement process.

In section 3 of this paper, we introduce a model system in \mathbb{R}^3 for a spin-1 quantum entity

subjected to Stern-Gerlach measurements, consisting of a joint system of two entangled¹ spin- $\frac{1}{2}$ quantum entities on which we introduce correlations of the second kind. In section 4, we study the states of this model system, and we'll find out that there are two kinds of states:

- i) If the system is in a *state of the first kind*, it behaves during the measurement as the joint system of two separated spin- $\frac{1}{2}$ quantum entities.
- ii) If the system is in a *state of the second kind*, it behaves during the measurement very similar as the joint system of two spin- $\frac{1}{2}$ quantum entities in a singlet state.

We consider a spin-1 entity that is prepared in a first Stern-Gerlach apparatus, and measured in a second one, and thus, the set of states that we consider is parameterized by directions in \mathbb{R}^3 , i.e., we consider a set of coherent spin-1 states. As been shown by Aaerge in [5], the hypothesis that the rotation invariant parameters of a spin-1 system are redundant “does not have dramatic consequences for the standard applications of quantum mechanics”.

2 Representation of a Spin- $\frac{1}{2}$ Quantum Entity

In this section, we represent the states of a spin- $\frac{1}{2}$ quantum entity on a sphere S .

2.1 The Transition Probability for Spin- $\frac{1}{2}$ Quantum Entity

The transition probability depends only on the relative position of the Stern-Gerlach apparatus in which we prepare the entity and a second Stern-Gerlach with which we measure the spin. We represent such a measurement by the Euler angles α, β, γ , and denote it as $e_{\alpha, \beta, \gamma}$. If the initial state corresponds with a spin quantum number $s = +\frac{1}{2}$ we denote it as p_+^0 , and if it corresponds with a spin quantum number $s = -\frac{1}{2}$ we denote it as p_-^0 . We represent p_+^0 by the vector $\psi_+^0 = (1, 0) \in \mathbb{C}^2$ and p_-^0 by $\psi_-^0 = (0, 1) \in \mathbb{C}^2$. The eigenstates corresponding to a measurement $e_{\alpha, \beta, \gamma}$ are the same as the ones we obtain when we rotate the initial states by an active rotation characterized by the Euler angles α, β, γ . This active rotation is represented by a unitary operator acting on \mathbb{C}^2 that corresponds with the following matrix (see [6] and [7]):

$$M_{\alpha, \beta, \gamma} = \begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos \frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin \frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin \frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos \frac{\beta}{2} \end{pmatrix} \quad (1)$$

Thus, for the measurement $e_{\alpha, \beta, \gamma}$ we have a set of eigenstates represented by the following vectors:

$$\psi_+^{\alpha, \beta, \gamma} = M_{\alpha, \beta, \gamma} \psi_+^0 = (e^{-i\frac{\alpha}{2}} \cos \frac{\beta}{2}, e^{i\frac{\alpha}{2}} \sin \frac{\beta}{2}) e^{-i\frac{\gamma}{2}} \quad (2a)$$

¹In this paper we call entities entangled whenever they are not separated.

$$\psi_{-}^{\alpha,\beta,\gamma} = M_{\alpha,\beta,\gamma} \psi_{-}^0 = (-e^{-i\frac{\alpha}{2}} \sin \frac{\beta}{2}, e^{i\frac{\alpha}{2}} \cos \frac{\beta}{2}) e^{i\frac{\gamma}{2}} \quad (2b)$$

The vectors in eq(2a) and eq(2b) that correspond to different values of γ (for fixed α and β) represent the same states. As a consequence, we omit the superscript γ in the notations for the vectors and the measurements. We represent the states corresponding to the vectors in eq(2a) and eq(2b) respectively by $p_{+}^{\alpha,\beta}$ and $p_{-}^{\alpha,\beta}$. Thus we have $p_{+}^{0,0} = p_{+}^0$ and $p_{-}^{0,0} = p_{-}^0$. We also have:

$$p_{+}^{\alpha,\beta} = p_{-}^{\alpha+\pi,\pi-\beta} \quad (3)$$

We denote the probability to obtain a state $p_{+}^{\alpha,\beta}$ in a measurement $e_{\alpha,\beta}$ on an entity in a state p_{+}^0 as $P_{+,+}^{\alpha,\beta}$, and the probability to obtain $p_{-}^{\alpha,\beta}$ in a measurement $e_{\alpha,\beta}$ on an entity in a state p_{+}^0 as $P_{+,-}^{\alpha,\beta}$. Analogously we define $P_{-,+}^{\alpha,\beta}$ and $P_{-,-}^{\alpha,\beta}$. We have:

$$P_{+,+}^{\alpha,\beta} = | \langle \psi_{+}^0 | \psi_{+}^{\alpha,\beta} \rangle |^2 = \cos^2 \frac{\beta}{2} = \frac{1 + \cos \beta}{2} \quad (4a)$$

$$P_{-,-}^{\alpha,\beta} = \frac{1 + \cos \beta}{2} \quad (4b)$$

$$P_{+,-}^{\alpha,\beta} = P_{-,+}^{\alpha,\beta} = \sin^2 \frac{\beta}{2} = \frac{1 - \cos \beta}{2} \quad (4c)$$

2.2 Representation of the Spin- $\frac{1}{2}$ States in S

The set of states of a spin- $\frac{1}{2}$ entity is given by (see eq.(3)):

$$\Sigma_{\frac{1}{2}} = \{ p_{+}^{\alpha,\beta} | \alpha \in [0, 2\pi], \beta \in [0, \pi] \} \quad (5)$$

Let S be a unit sphere in \mathbb{R}^3 with its center in the origin. We represent every state $p_{+}^{\alpha,\beta} \in \Sigma_{\frac{1}{2}}$ by the point in S with coordinates $(\cos \alpha \sin \beta, \sin \alpha \sin \beta, \cos \beta)$. It is clear (as a consequence of the definition of the Euler angles), that the representation of $\Sigma_{\frac{1}{2}}$ in S is one to one and onto.

3 A Model System for a Spin-1 Quantum Entity.

In this section we'll introduce the model system for a spin-1 quantum entity. First we represent the states of a spin-1 quantum entity in $S \times S$, i.e., we represent every spin state as two points on a sphere².

²The representation introduced in this paper is not the same as the Majorana representation for all spin-1 states, presented in [8]. One can easily verify (by applying the construction for the Majorana representation presented in [9]) that our representation could not be extended in a continuous way to a one to one representation of all spin-1 states.

3.1 The Transition Probability for Coherent Spin-1 States

We proceed along the same lines as in section 2.1. We denote a measurement characterized by the Euler angles α, β, γ as $e_{\alpha, \beta, \gamma}$. If the initial state of the entity corresponds with a spin quantum number $s = +1$, we denote it as p_+^0 . If $s = 0$, we denote it as p_0^0 and if $s = -1$ as p_-^0 . We represent p_+^0 by the vector $\psi_+^0 = (1, 0, 0) \in \mathbb{C}^3$, p_0^0 by $\psi_0^0 = (0, 1, 0) \in \mathbb{C}^3$ and p_-^0 by $\psi_-^0 = (0, 0, 1) \in \mathbb{C}^3$. The eigenstates corresponding to a measurement $e_{\alpha, \beta, \gamma}$ are the same as the ones we obtain when we rotate the initial states by an active rotation characterized by the Euler angles α, β, γ . This active rotation is represented by a unitary operator acting on \mathbb{C}^3 following matrix (see [6] and [7]):

$$M_{\alpha, \beta, \gamma} = \begin{pmatrix} \frac{1+\cos\beta}{2}e^{-i\frac{\alpha+\gamma}{2}} & -\frac{\sin\beta}{\sqrt{2}}e^{-i\frac{\alpha}{2}} & \frac{1-\cos\beta}{2}e^{-i\frac{\alpha-\gamma}{2}} \\ \frac{\sin\beta}{\sqrt{2}}e^{-i\frac{\gamma}{2}} & \cos\beta & -\frac{\sin\beta}{\sqrt{2}}e^{i\frac{\gamma}{2}} \\ \frac{1-\cos\beta}{2}e^{i\frac{\alpha-\gamma}{2}} & \frac{\sin\beta}{\sqrt{2}}e^{i\frac{\alpha}{2}} & \frac{1+\cos\beta}{2}e^{i\frac{\alpha+\gamma}{2}} \end{pmatrix} \quad (6)$$

Thus, for the measurement $e_{\alpha, \beta, \gamma}$ we have a set of eigenstates represented by the following vectors:

$$\psi_+^{\alpha, \beta, \gamma} = M_{\alpha, \beta, \gamma} \psi_+^0 = \left(\frac{1+\cos\beta}{2}e^{-i\frac{\alpha}{2}}, \frac{\sin\beta}{\sqrt{2}}, \frac{1-\cos\beta}{2}e^{i\frac{\alpha}{2}} \right) e^{-i\frac{\gamma}{2}} \quad (7a)$$

$$\psi_0^{\alpha, \beta, \gamma} = M_{\alpha, \beta, \gamma} \psi_0^0 = \left(-\frac{\sin\beta}{\sqrt{2}}e^{-i\frac{\alpha}{2}}, \cos\beta, \frac{\sin\beta}{\sqrt{2}}e^{i\frac{\alpha}{2}} \right) \quad (7b)$$

$$\psi_-^{\alpha, \beta, \gamma} = M_{\alpha, \beta, \gamma} \psi_-^0 = \left(\frac{1-\cos\beta}{2}e^{-i\frac{\alpha}{2}}, -\frac{\sin\beta}{\sqrt{2}}, \frac{1+\cos\beta}{2}e^{i\frac{\alpha}{2}} \right) e^{i\frac{\gamma}{2}} \quad (7c)$$

As been motivated in section 2.1, we represent the states corresponding to the vectors in eq(7a), eq(7b) and eq(7c), respectively by $p_+^{\alpha, \beta}$, $p_0^{\alpha, \beta}$ and $p_-^{\alpha, \beta}$, and we omit the subscript γ in the notation of the measurement. We have:

$$p_+^{\alpha, \beta} = p_-^{\alpha, \pi+\beta} \quad (8)$$

$$p_0^{\alpha, \beta} = p_0^{\alpha, \pi+\beta} \quad (9)$$

Using similar notations as introduced in section 2.1, we find the transition probabilities through square of the Hilbert in-product on \mathbb{C}^3 :

$$P_{0,0}^{\alpha, \beta} = | \langle \psi_0^0 | \psi_0^{\alpha, \beta} \rangle |^2 = \cos^2 \beta \quad (10a)$$

$$P_{+,+}^{\alpha, \beta} = P_{-,-}^{\alpha, \beta} = \left(\frac{1+\cos\beta}{2} \right)^2 \quad (10b)$$

$$P_{+,-}^{\alpha, \beta} = P_{-,+}^{\alpha, \beta} = \left(\frac{1-\cos\beta}{2} \right)^2 \quad (10c)$$

$$P_{+,0}^{\alpha, \beta} = P_{0,+}^{\alpha, \beta} = P_{-,0}^{\alpha, \beta} = P_{0,-}^{\alpha, \beta} = \frac{\sin^2 \beta}{2} \quad (10d)$$

Again the transition probabilities depend only on the angle β .

3.2 Representation of the Coherent Spin-1 States in $S \times S$

We define the following two subsets of the set of coherent spin states:

$$\Sigma_+ = \{p_+^{\alpha,\beta} | \alpha \in [0, 2\pi], \beta \in [0, \pi]\} \quad (11a)$$

$$\Sigma_0 = \{p_0^{\alpha,\beta} | \alpha \in [0, 2\pi], \beta \in [0, \pi]\} \quad (11b)$$

We call the states contained in Σ_+ states of the first kind, the states contained in Σ_0 states of the second kind. The set of all coherent spin-1 states is (see eq(8), eq(9), eq(11a) and eq(11b)):

$$\Sigma_1 = \Sigma_+ \cup \Sigma_0 \quad (12)$$

If $p_+^{\alpha,\beta} \in \Sigma_+$ we represent this state by two identical points in S (one point of $S \times S$) with coordinates $(\cos\alpha\sin\beta, \sin\alpha\sin\beta, \cos\beta)$. If $p_0^{\alpha,\beta} \in \Sigma_0$ we represent this state by two points in S with respective coordinates $(\cos\alpha\sin\beta, \sin\alpha\sin\beta, \cos\beta)$ and $(-\cos\alpha\sin\beta, -\sin\alpha\sin\beta, -\cos\beta)$. It is clear that this representation of Σ_1 in $S \times S$ is one to one (see eq.(9)).

3.3 The Model System

In this subsection we introduce a model system. Then we prove that this model system is a representation for a spin-1 quantum entity.

Suppose that the states of model system can be represented by the same subset of $S \times S$, as we used in the previous section for the representation of the coherent states of a spin-1 quantum entity. A state represented as two identical points on the sphere with coordinates v will be denoted as $p_{v,v}$, and a state represented as two diametrically opposite points with coordinates v and $-v$ as $p_{v,-v}$. We choose a set of coordinates such $v = (0, 0, 1)$. We define a measurement e_u on the system in a state $p_{v,v}$ in the following way:

- i) We consider the system as a joint system of two separated spin- $\frac{1}{2}$ quantum entities in a state p_v .
- ii) On both entities perform a measurement³ with eigenstates p_u and p_{-u} .

We define the measurement e_u on the system in a state $p_{-v,v}$ in the following way:

- i) We consider the system as a joint system of two spin- $\frac{1}{2}$ quantum entities that are entangled, one of them in a state p_v , the other in a state p_{-v} .
- ii) On one of these two entities, that we denote as \mathcal{S}_1 , we perform a measurement with eigenstates p_u and p_{-u} . Let $u = (u_1, u_2, u_3)$. If, as a consequence of this measurement, we obtain a state p_u for \mathcal{S}_1 , then the state of the other entity (denoted as \mathcal{S}_2) changes to $p_{v'}$, where $v' = (u_1, u_2, -u_3)$. If we obtain a state p_{-u} for \mathcal{S}_1 , then the state of \mathcal{S}_2 changes to $p_{v'}$.

³We consider measurements on spin- $\frac{1}{2}$ quantum entities as presented in section 2.

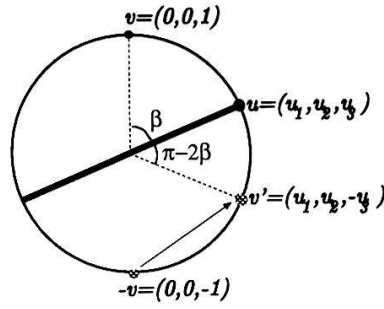


Figure 1: Illustration of the e_u measurements on the model system in a state $p_{v,-v}$. In this illustration, the black spots represent position of \mathcal{S}_1 .

iii) We perform a measurement with eigenstates p_u and p_{-u} on \mathcal{S}_2 .

There are three possible outcome states for e_u : $p_{u,u}$, $p_{-u,-u}$ or $p_{-u,u}$. Denote by β the angle between the vectors u and v . We can calculate the transition probabilities. Denote the probability to become a state $p_{u,u'}$ in the measurement e_u on the system in a state $p_{v,v'}$ as $P_{v,v',u,u'}$. We have: $P_{v,v,u,u} = (\frac{1+\cos\beta}{2})^2$, $P_{v,v,-u,-u} = (\frac{1-\cos\beta}{2})^2$ and $P_{v,v,u,-u} = \frac{1+\cos\beta}{2} \cdot \frac{1-\cos\beta}{2} + \frac{1-\cos\beta}{2} \cdot \frac{1+\cos\beta}{2} = \frac{\sin^2\beta}{2}$. Following section 3.2, we can identify the state $p_{v,v}$ with the state $p_+^0 \in \Sigma_1$ of a spin-1 quantum entity, and the state $p_{u,u}$ with the state $p_+^{\alpha,\beta} \in \Sigma_1$ and the state $p_{u,-u}$ with the state $p_o^{\alpha,\beta} \in \Sigma_1$. According to eq.(10b), eq.(10c) and eq.(10d), we find the same transition probabilities for the model system as for a spin-1 quantum entity, for an initial state contained in Σ_+ .

For an initial state $p_{v,-v}$ we have to take the two possibilities into account: \mathcal{S}_1 is in a state p_v or \mathcal{S}_2 is in a state p_v . For example: the probability to become a state $p_{u,u}$ in the measurement e_u on the entity in a state $p_{v,-v}$, if \mathcal{S}_1 is in a state p_v , is equal to the the probability to obtain the state p_u in the measurement on \mathcal{S}_1 , multiplied by the probability to obtain a state p_u in the measurement on \mathcal{S}_2 . The probability to obtain p_u in this second measurement is $\frac{1}{2}(1 + \cos[\pi - 2\beta])$. We have:

$$P_{v,-v,u,u} = \frac{1}{2} \cdot \frac{1+\cos\beta}{2} \cdot \frac{1+\cos[\pi-2\beta]}{2} + \frac{1}{2} \cdot \frac{1-\cos\beta}{2} \cdot \frac{1+\cos[\pi-2\beta]}{2} = \frac{1+\cos[\pi-2\beta]}{4} = \frac{\sin^2\beta}{2}$$

$$\begin{aligned} P_{v,-v,u,-u} &= \frac{1}{2} \cdot \frac{1+\cos\beta}{2} \cdot \frac{1-\cos[\pi-2\beta]}{2} + \frac{1}{2} \cdot \frac{1-\cos\beta}{2} \cdot \frac{1-\cos[\pi-2\beta]}{2} \\ &\quad + \frac{1}{2} \cdot \frac{1-\cos\beta}{2} \cdot \frac{1-\cos[\pi-2\beta]}{2} + \frac{1}{2} \cdot \frac{1+\cos\beta}{2} \cdot \frac{1-\cos[\pi-2\beta]}{2} = \frac{1-\cos[\pi-2\beta]}{2} = \cos^2\beta \end{aligned}$$

and

$$P_{v,-v,-u,-u} = \frac{1}{2} \cdot \frac{1-\cos\beta}{2} \cdot \frac{1+\cos[\pi-2\beta]}{2} + \frac{1}{2} \cdot \frac{1+\cos\beta}{2} \cdot \frac{1+\cos[\pi-2\beta]}{2} = \frac{1+\cos[\pi-2\beta]}{4} = \frac{\sin^2\beta}{2}$$

Thus, if we identify $p_{v,-v}$ with the state $p_+^0 \in \Sigma_1$ of a spin-1 quantum entity, we can compare eq.(10a) and eq.(10d) with $P_{v,-v,u,u}$, $P_{v,-v,u,-u}$ and $P_{v,-v,-u,-u}$. Again we find the same transition probabilities for the model system as for a spin-1 quantum entity. Thus, this model system is a representation for a spin-1 quantum entity with Σ_1 as set of states.

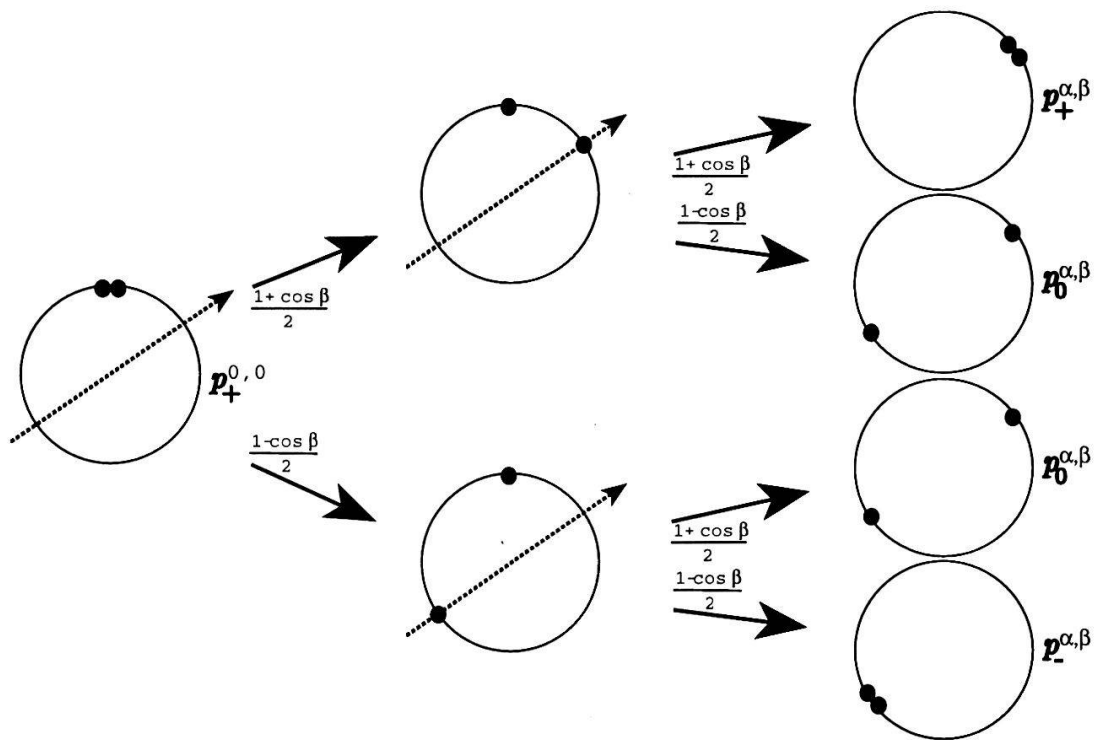


Figure 2: Illustration going with the calculation of the transition probabilities of the model system for spin-1 for an initial state in Σ_1 .

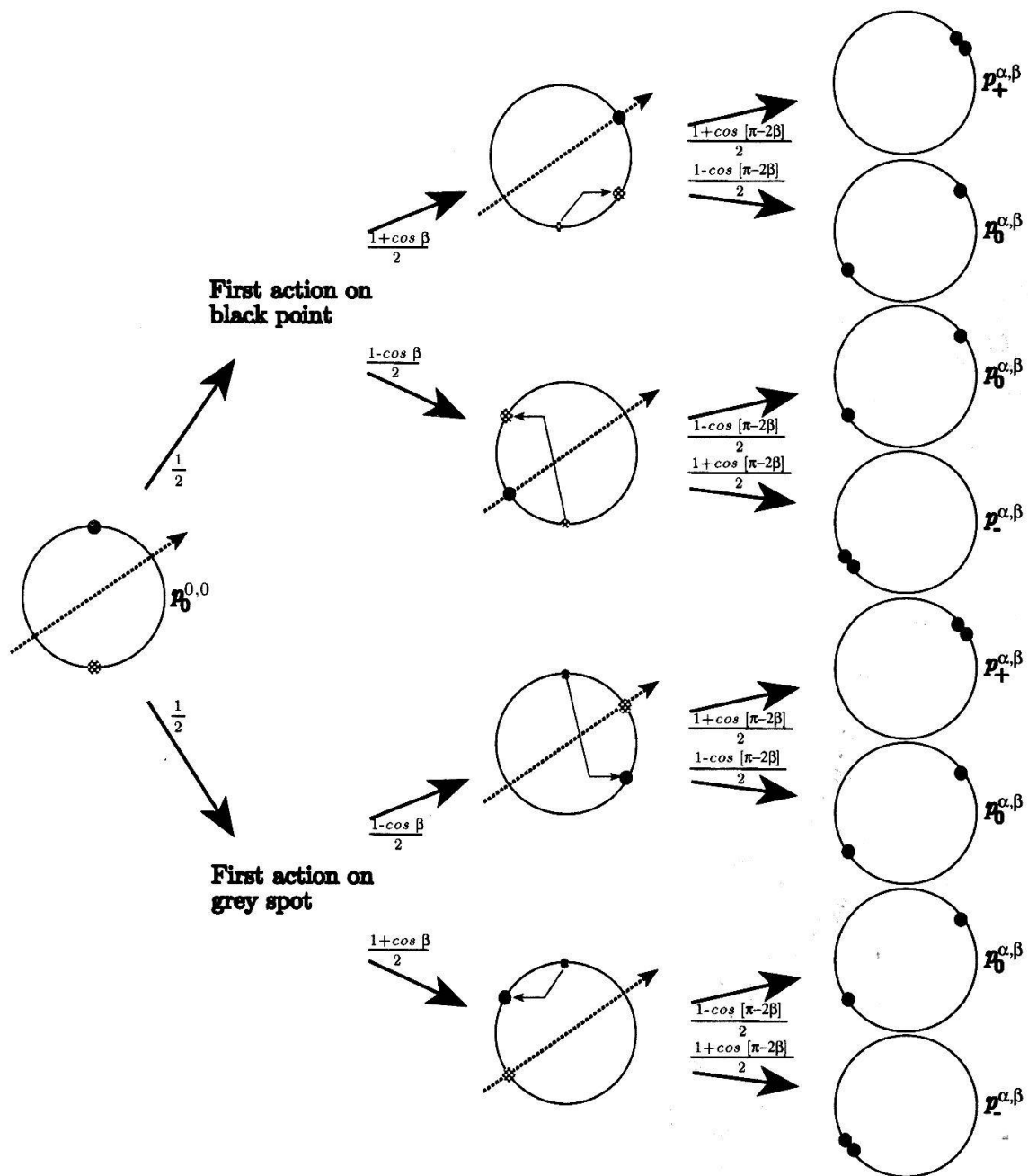


Figure 3: Illustration going with the calculation of the transition probabilities of the model system for spin-1, for an initial state in Σ_0 .

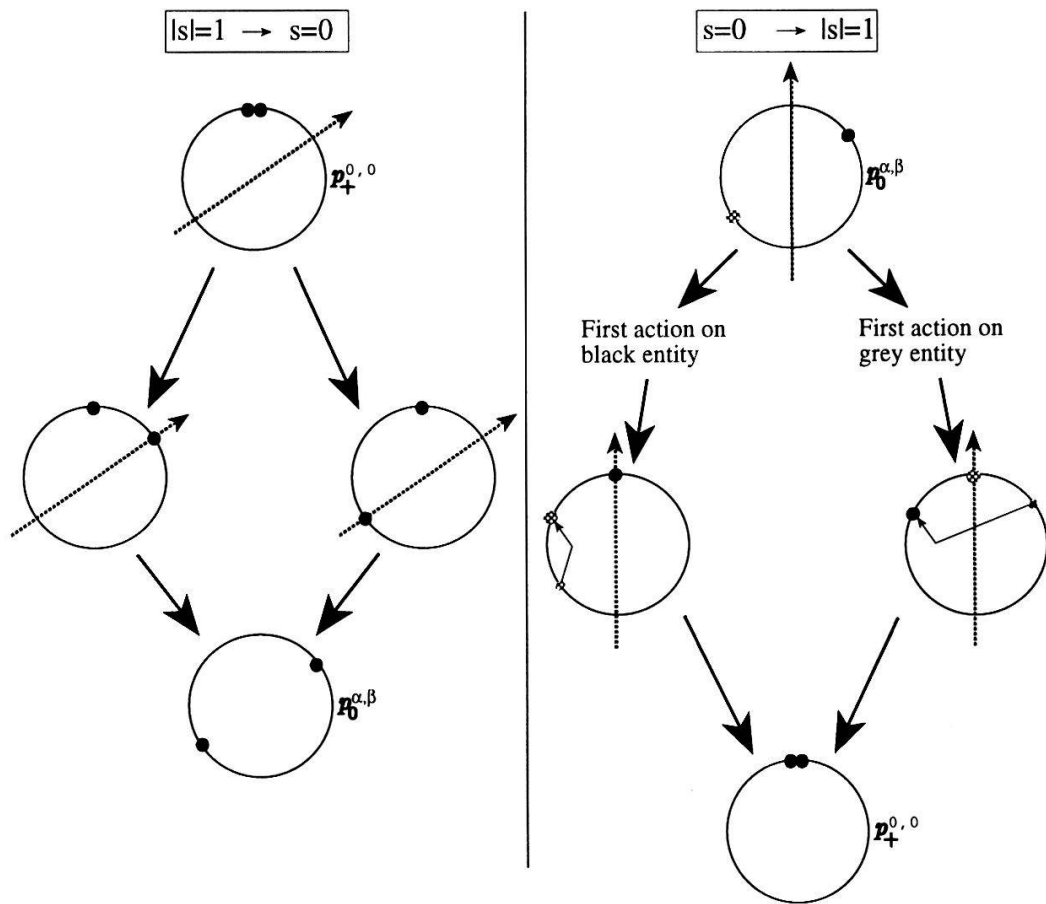


Figure 4: Two state transitions corresponding with two completely different mechanisms, but with the same probability of appearance. The symmetry of the transition probability does not necessarily imply a process that is reversible from a geometrical point of view.

4 The States of the Model System

As we saw in section 3.3, Σ_+ corresponds to the states of a joint system of two separated spin- $\frac{1}{2}$ quantum entities, and thus, *a spin-1 quantum entity in a state in Σ_+ behaves as two separated spin- $\frac{1}{2}$ quantum entities.*

In our model system, the correlations are created during the measurement, and thus, *they are of the second kind.* In [4], Aerts shows that due to the presence of correlations of the second kind, the singlet state is a new state, not contained in $\mathcal{H}(\mathbb{C}^2) \times \mathcal{H}(\mathbb{C}^2)$ ($\mathcal{H}(X)$ are the rays of X), but in $\mathcal{H}(\mathbb{C}^2 \otimes \mathbb{C}^2)$. Thus, it are these kinds of correlations that are responsible for the 'quantum nature' of the singlet state. All this indicates a different approach towards the description of joint systems: *if a joint system consists of two separated entities, the states are represented by the cartesian product of the state spaces of these entities; if due to the interaction between the measurement apparatus and the joint system, correlations of the second kind are created, new states occur.* The states in our model system, are also new states: we know that they can be represented in $\mathcal{H}(\mathbb{C}^3)$, a subspace of $\mathcal{P}(\mathbb{C}^2 \otimes \mathbb{C}^2)$, and one can verify that, although we have represented the initial states of the model system in $S \times S$, these states do not correspond to the subset $\mathcal{H}(\mathbb{C}^2) \times \mathcal{P}(\mathbb{C}^2)$ of $\mathcal{H}(\mathbb{C}^2 \otimes \mathbb{C}^2)$.

The symmetry of the transition probability of the presented model system seems to be 'accidental', in the sense that the two transitions $p_+^{0,0} \rightarrow p_0^{\alpha,\beta}$ (or $p_+^{\alpha,\beta} \rightarrow p_0^{0,0}$, or $p_-^{0,0} \rightarrow p_0^{\alpha,\beta}$) and $p_0^{\alpha,\beta} \rightarrow p_+^{0,0}$ (or $p_0^{0,0} \rightarrow p_+^{\alpha,\beta}$, or $p_0^{\alpha,\beta} \rightarrow p_-^{0,0}$) correspond to processes with a completely different geometry. The geometry of the transition $p_+^{0,0} \rightarrow p_0^{\alpha,\beta}$ corresponds to the one of a measurement on a joint system of two separated spin- $\frac{1}{2}$ quantum entities. The transition $p_0^{\alpha,\beta} \rightarrow p_+^{0,0}$ corresponds to a measurement on a joint system of two entangled entities.

5 Conclusion

In this paper, we showed that a spin-1 quantum entity can be interpreted as a joint system of two entangled spin- $\frac{1}{2}$ quantum entities. Moreover, by introducing correlations of the second kind on this joint system of two spin- $\frac{1}{2}$ quantum entities, we find states in $\mathcal{H}(\mathbb{C}^2 \otimes \mathbb{C}^2)$ that are not contained in $\mathcal{H}(\mathbb{C}^2) \times \mathcal{H}(\mathbb{C}^2)$. Thus, this model system indicates a new way of looking at a spin-1 quantum entity, and motivates a new approach towards the description of joint systems.

6 Acknowledgements

We thank Prof. D. Aerts for the many discussions concerning this topic. We thank Prof. C. Piron for pointing out references related to this topic. This work is supported by the 'Federale Diensten voor Wetenschappelijke, Technische en Culturele Aangelegenheden' in the

framework of IUAP-III n°9.

References

- [1] D. Aerts, *The one and the many*, (Doctoral Dissertation, Free University of Brussels 1981)
- [2] C. Piron, *Foundations of Quantum Physics*, (Benjamin Reading 1975)
- [3] D. Aerts, *Found. Phys.* **24**, 1227 (1994)
- [4] D. Aerts, *Helv. Phys. Acta.* **64**, 1 (1991)
- [5] T. Aaberge, *Helv. Phys. Acta.* **67**, 127 (1994)
- [6] A. Perelomov, *Generalized Coherent States and Their Applications* (Springer 1986)
- [7] E. P. Wigner, *Group Theory and its Applications to the Quantum Mechanics of Atomic Spectra*, (Academic press 1959)
- [8] E. Majorana, *Nuovo Cimento* **9**, 43 (1932)
- [9] H. Bacry, *J. Math. Phys.* **15**, 1686 (1974)