

**Zeitschrift:** Helvetica Physica Acta  
**Band:** 66 (1993)  
**Heft:** 2  
  
**Artikel:** Distortion analyticity for N-particle hamiltonians  
**Autor:** Gérard, C.  
**DOI:** <https://doi.org/10.5169/seals-116569>

### **Nutzungsbedingungen**

Die ETH-Bibliothek ist die Anbieterin der digitalisierten Zeitschriften auf E-Periodica. Sie besitzt keine Urheberrechte an den Zeitschriften und ist nicht verantwortlich für deren Inhalte. Die Rechte liegen in der Regel bei den Herausgebern beziehungsweise den externen Rechteinhabern. Das Veröffentlichen von Bildern in Print- und Online-Publikationen sowie auf Social Media-Kanälen oder Webseiten ist nur mit vorheriger Genehmigung der Rechteinhaber erlaubt. [Mehr erfahren](#)

### **Conditions d'utilisation**

L'ETH Library est le fournisseur des revues numérisées. Elle ne détient aucun droit d'auteur sur les revues et n'est pas responsable de leur contenu. En règle générale, les droits sont détenus par les éditeurs ou les détenteurs de droits externes. La reproduction d'images dans des publications imprimées ou en ligne ainsi que sur des canaux de médias sociaux ou des sites web n'est autorisée qu'avec l'accord préalable des détenteurs des droits. [En savoir plus](#)

### **Terms of use**

The ETH Library is the provider of the digitised journals. It does not own any copyrights to the journals and is not responsible for their content. The rights usually lie with the publishers or the external rights holders. Publishing images in print and online publications, as well as on social media channels or websites, is only permitted with the prior consent of the rights holders. [Find out more](#)

**Download PDF:** 13.09.2025

**ETH-Bibliothek Zürich, E-Periodica, <https://www.e-periodica.ch>**

# Distortion analyticity for $N$ -particle Hamiltonians

C.Gérard \*

Institute for Advanced Study  
Olden Lane Princeton NJ 08540 USA

October 1992

(28. X. 1992)

## Abstract

We define resonances for  $N$ -particle Hamiltonians with pair potentials which are not dilation analytic.

## 1 Introduction

We consider in this paper  $N$ -particle Hamiltonians

$$H = \sum_{i=1}^N \frac{1}{2m_i} \Delta_{x_i} + \sum_{i<j} V_{ij}(x_i - x_j)$$

on  $L^2(\mathbf{R}^{N\nu})$ . A rigorous mathematical theory of resonances for such Hamiltonians has been developed in the pioneering papers by Aguilar-Combes [Ag-Co] for 2-particle Hamiltonians and Balslev-Combes [Ba-Co] for arbitrary  $N$ . In this approach resonances are defined as complex eigenvalues of some non self-adjoint deformation of  $H$ . However in these two papers it is essential that the pair potentials are *dilation analytic*, which roughly means that the  $V_{ij}(y)$  have to extend holomorphically in a cone  $\{|Imz| \leq C|Rez|\}$ . For some applications (like for example  $N$ -particle Hamiltonians where some particles have infinite mass) it proved necessary to extend the class of potentials in order to accept potentials which are analytic only near infinity. This was done for 2-particle Hamiltonians by a number of authors (see [Si], [S], [Cy], [Hu]), by introducing variants of the dilation method.

However it seems that no such results are known for  $N$ -particle Hamiltonians when  $N \geq 3$ . For example there is no definition of resonances in the literature even when the pair potentials have compact support.

---

\*Permanent adress: Centre de Mathématiques, Ecole Polytechnique 91128 Palaiseau Cedex France.

The goal of this note is to fill this gap by defining resonances for  $N$ -particle Hamiltonians when the pair potentials are analytic only near infinity. This is done using the analytic distortion method of Hunziker [Hu] with a vector field which respects the  $N$ -particle structure of the potential. This vector has been originally introduced by Graf [Gr] to prove propagation estimates for  $N$ -particle Hamiltonians and is now a fundamental tool in the scattering theory for such systems.

Let us now describe more in details the class of Hamiltonians which we will consider. We study a straightforward extension of  $N$ -body Hamiltonians called *Agmon Hamiltonians* (cf [Ag]). One considers a finite dimensional real vector space  $X$  with a positive definite quadratic form  $g(x, x)$ , and a finite family  $\{X_a\}$ ,  $a \in A$  of linear vector subspaces of  $X$  which is closed under intersection and obeying  $\bigcap_{a \in A} X_a = \{0\}$  and  $X \in \{X_a\}$ . One denotes by  $X^a$  the space  $X_a^\perp$ , by  $\pi^a$ ,  $\pi_a$  the orthogonal projections on  $X^a$  and  $X_a$ .

On  $A$  one puts a partial ordering by saying that  $b \leq a$  if  $X^b \subset X^a$ . With this ordering  $A$  is a lattice and one gets that  $X_{a_{\max}} = \{0\}$  and  $X_{a_{\min}} = X$ . Let  $D_x = \frac{1}{i} \partial_x$  and let  $\langle x \rangle = (1 + g(x, x))^{1/2}$ . For  $a \in A$ , one denotes by  $\#a$  the maximal number  $k$  such that  $a_1 = a < a_2 < \dots < a_k = a_{\max}$ .

If  $N = \#a_{\min}$ , one defines a (generalized)  $N$ -body Hamiltonian by :

$$H = \frac{1}{2} \tilde{g}(D_x, D_x) + V(x),$$

where :  $V(x) = \sum_{a \in A} V_a(\pi^a x)$  and  $\tilde{g}$  is the dual quadratic form on  $X'$  associated with  $g$ . For simplicity of notations, we will simply denote  $\tilde{g}(D_x, D_x)$  by  $D_x^2$ .

For  $a \in A$ , we denote by  $H_a$  the Hamiltonian  $H - I_a(x)$ , where  $I_a(x) = \sum_{b \not\leq a} V_b(x^b)$ . One has also  $H_a = \frac{1}{2} D_a^2 + H^a$ , where  $H^a$  is the Hamiltonian acting on  $L^2(X^a)$  defined by  $H^a = \frac{1}{2} D^{a2} + V^a(x^a)$  for  $V^a(x^a) = \sum_{b \leq a} V_b(x^b)$ .

We will assume that the potentials  $V_a$  satisfy the following hypotheses :

$$H1) V_a \in L^\infty(X^a).$$

$H2) V_a$  extends holomorphically in

$$\{z \in \mathbb{C}^{n^*} \mid |Re z| \leq R, |Im z| < \epsilon_0 |Re z|\},$$

for some  $R, \epsilon_0$ , and satisfies in this region :

$$H3) \lim_{z \rightarrow \infty} V_a(z) = 0.$$

The condition that  $V_a \in L^\infty(X^a)$  is purely for illustrative purposes. The extension to singular potentials is easy.

Let us now give the plan of this paper. In Section 2, we recall the definition of the Graf's vector field and prove two important properties. In Section 3 we define resonances as eigenvalues of the distorted Hamiltonian and prove that they coincide with poles of the meromorphic continuation of the resolvent.

## 2 The distortion vector field

To define the complex distortion, we will use a vector field originally introduced by Graf [Gr] to prove propagation estimates for  $N$ -body Hamiltonians. For the reader's convenience, we will briefly recall its construction.

Let us first introduce some notations. For  $a \leq b$  one defines :

$$x_a^b := \pi^b x_a = x^b - x^a = \pi_a x^b.$$

Note that :

$$(x_a^b)^2 = x_a^2 - x_b^2 = x^{b2} - x^{a2}.$$

One puts then :

$$(2.1) \quad J_a(x) := \prod_{a < f} F((x_a^f)^2 > q_a^f) \prod_{g < a} F((x_g^a)^2 \leq q_g^a),$$

where  $F(x \in A)$  denotes the characteristic function of  $A$ . The constants  $q_b^a$  are chosen equal to :

$$q_b^a = q^a - q^b,$$

where

$$q^a := \begin{cases} q^{a-1} & \text{if } a \neq a_{\min} \\ 0 & \text{if } a = a_{\min}. \end{cases}$$

For a mollifier  $\phi \in C_0^\infty(X)$  with :

$$\phi \geq 0, \int \phi(x) dx = 1, \int x \phi(x) dx = 0, \text{supp } \phi \subset \{|x| \leq \sigma\},$$

one then defines :

$$j_a(x) := J_a \star \phi(x).$$

We will use the following properties of  $j_a$  (see [Gr]) :

**Lemma 2.1**

- i)  $\sum_{a \in \mathcal{A}} j_a(x) = 1.$
- ii)  $\exists C_0$  such that  $|x^a| \leq C_0$  on  $\text{supp } j_a.$
- iii)  $\exists C_1$  such that  $\forall b \not\leq a, |x^b| \geq C_1$ , on  $\text{supp } j_a.$

We then define the distortion vector field :

$$v_C(x) := \sum_{a \in \mathcal{A}} j_a\left(\frac{x}{C}\right) x_a.$$

the constant  $C$  will have to be chosen large enough later. For ease of notations we will usually forget the subscript  $C$  and write simply  $v(x)$  for  $v_C(x)$ . We remark that if  $N = 2$ , then  $V_C(x)$  is identical to the distortion vector field of [Hu]. Note the following estimate, which follows directly from Lemma 2.1 :

$$(2.2) \quad |\partial_x^\alpha v_C(x)| \leq C_\alpha, \forall \alpha \in \mathbb{N}.$$

Let us now recall the definition of the *distortion* associated with  $v_C$  given in [Hu]. Since  $\nabla_x v_C(x) = O(1)$ , the mapping

$$X \ni x \mapsto x + \theta v_C(x)$$

is invertible for  $\theta \in \mathbb{R}$ ,  $|\theta| \leq c_0$ . So we can define the unitary transformation  $U_\theta$  on  $L^2(X)$  by

$$U_\theta u(x) := J_\theta^{\frac{1}{2}} u(x + \theta v_C(x)),$$

for  $J_\theta = \det(\delta_{ik} + \partial_k v^i)$ . One then puts

$$H_\theta := U_\theta H U_\theta^{-1} = \frac{1}{2} U_\theta D_x^2 U_\theta^{-1} + V_\theta(x).$$

Here

$$U_\theta D_x^2 U_\theta^{-1} = J_\theta^{-\frac{1}{2}} D_i J_\theta^{ik} J_\theta J_\theta^{mk} D_k J_\theta^{-\frac{1}{2}},$$

$$V_\theta(x) = V(x + \theta v_C(x)),$$

where  $(J_\theta^{ik})$  is the inverse of the matrix  $(J_\theta)_{ik}$  and summation over repeated indices is understood.

The goal is now to extend  $H_\theta$  to complex values of  $\theta$ . The first important property of  $v$  is the following :

**Proposition 2.2** *Assume that :*

$$V(x) = \sum_{a \in \mathcal{A}} V_a(x^a),$$

where  $V_a$  satisfy the conditions  $H$ ). Then for  $C$  large enough, there exist  $c_0, c_1$  such that :

$$\theta \mapsto V_\theta(x) := V(x + \theta v_C(x))$$

extends from  $\theta \in \mathbb{R}$  as a function holomorphic in

$$\{\theta \in \mathbb{C} \mid |\operatorname{Re} \theta| \leq c_0, |\operatorname{Im} \theta| \leq c_1\}$$

with values in  $L^\infty(X)$ .

*proof.* it suffices to consider a potential  $V_b(x^b)$ . For  $\theta \in \mathbf{R}$ , we write :

$$\begin{aligned} V_b(x^b + \theta \pi^b v(x)) &= V_b(x^b + \theta \sum_{a \in \mathcal{A}} j_a(\frac{x}{C}) \pi^b x_a) = \\ &V_b(x^b + \theta \sum_{a \in \mathcal{A}, b \not\leq a} j_a(\frac{x}{C}) \pi^b x_a), \end{aligned}$$

since  $\pi^b x_a = 0$  if  $b \leq a$ . Next we observe that on  $\text{supp} j_a(\frac{x}{C})$ , one has :

$$(2.3) \quad |x^a| \leq C_0 C, \quad |x^b| \geq C_1 C, \quad \forall b \not\leq a.$$

So :

$$(2.4) \quad |\pi^b x_a| = |\pi^b x - \pi^b x_a| \leq |x^b| + C_0 C.$$

By (2.3), we see that

$$\sum_{a \in \mathcal{A}, b \not\leq a} j_a(\frac{x}{C}) \pi^b x_a$$

is supported in  $\{|x^b| \geq C_1 C\}$ , so that

$$V_b(x^b + \theta \sum_{a \in \mathcal{A}, b \not\leq a} j_a(\frac{x}{C}) \pi^b x_a)$$

is clearly analytic in  $\theta$  for  $|x^b| \leq C_1 C$ . By  $H$ ), we can now pick  $C$  large enough such that if  $|x^b| \geq C_1 C$  then

$$\theta \mapsto V_b(x^b + \theta y^b)$$

is holomorphic in  $|\text{Im} \theta| \leq c_0$  for some  $c_0 > 0$ , uniformly for  $|y^b| \leq c_2 |x^b|$ . But using (2.4), we have :

$$| \sum_{a \in \mathcal{A}, b \not\leq a} j_a(\frac{x}{C}) \pi^b x_a | \leq c_2 |x^b|,$$

in  $\{|x^b| \geq C_1 C\}$ , so that

$$\theta \mapsto V_b(x^b + \theta \pi^b v(x))$$

is holomorphic in  $\theta$  as claimed.  $\square$

To state the second property, we introduce another partition of unity. It is well known (see for example [C.F.K.S]), that there exist a partition of unity :

$$(2.5) \quad 1 = \sum_{|a| \leq 2} q_a(x),$$

with the following properties :

$$q_{a_{\max}} \in C_0^\infty(|x| \leq 2),$$

$$\text{for } |a| = 2, \text{ } \text{supp} q_a \subset \{x \in X \mid |x| \geq 1, |x^b| \geq \epsilon_0 |x|, \forall b \not\leq a\},$$

$$\text{for } |a| = 2, \text{ } q_a \in C^\infty(X), \quad |\nabla_x q_a| \leq C \langle x \rangle^{-1}.$$

We will denote by  $q_{a,R}(x)$  the scaled functions  $q_a(\frac{x}{R})$ . In the next proposition, we will denote by  $v_C^a(x^a)$  a vector field on  $X^a$  defined exactly as  $v_C$ , replacing  $X$  by  $X^a$  and the set of indices  $\mathcal{A}$  by the subset  $\{b \in \mathcal{A} \mid b \leq a\}$ . Accordingly in the definition of the constants  $q_b^d$  for  $b, d \leq a$ , one has to replace  $\#b$  by  $\#^a b$  defined as the maximal number  $k$  such that :

$$b_1 = b < a_2 \cdots < b_k = a.$$

Note that by the Jordan-Dedekind chain condition, one has :

$$(2.6) \quad \#b + 1 = \#^a b + \#a.$$

**Proposition 2.3** *For any  $C > 0$ , there exist  $R > 0$  such that  $\forall a$  with  $\#a = 2$  one has :*

$$v_C(x) = x_a + v_C^a(x^a)$$

on  $\text{supp} q_{a,R}$ .

*proof.* let us consider

$$v_C(x) = \sum_b j_b\left(\frac{x}{C}\right)x_b.$$

On  $\text{supp} q_{a,R}$  one has

$$|x| \geq R, |x^b| \geq \epsilon_0 |x|, \text{ if } b \not\leq a.$$

So for  $R$  large enough, one has :

$$(2.7) \quad \begin{aligned} v_C(x) &= \sum_{b \leq a} j_b\left(\frac{x}{C}\right)x_b = \\ &\sum_{b \leq a} j_b\left(\frac{x}{C}\right)x_b^a + \\ &\left(\sum_{b \leq a} j_b\left(\frac{x}{C}\right)\right)x_a = \\ &\sum_{b \leq a} j_b\left(\frac{x}{C}\right)x_b^a + x_a \end{aligned}$$

on  $\text{supp} q_{a,R}$ . For  $b \leq a$ , we write :

$$j_b\left(\frac{x}{C}\right) = \int J_b\left(\frac{x}{C} + y\right)\phi(y)dy$$

and replace  $J_b$  by its expression given in (2.1). If  $x \in \text{supp} q_{a,R}$ ,  $y \in \text{supp} \phi$ , and  $f \not\leq a$ , one has :

$$\begin{aligned} \left(\frac{x_b^f}{C} + y_b^f\right)^2 &\geq \frac{1}{2} \frac{(\epsilon_0 R)^2}{C^2} - \frac{1}{2} \sigma^2 - \left(\frac{x^b}{C} + y^b\right)^2 \geq \\ &\frac{1}{2} \frac{(\epsilon_0 R)^2}{C^2} - \frac{1}{2} \sigma^2 - C_0^2, \end{aligned}$$

since  $|x^f| \geq \epsilon_0 R$  and  $|\frac{x^b}{C} + y^b| \leq C_0$  if  $J_b(\frac{x}{C} + y) \neq 0$ . If we put

$$J_b^a(x^a) := \Pi_{b < f \leq a} F((x_b^f)^2 > q_b^f) \Pi_{g < b} F((x_g^b)^2 \leq q_g^b),$$

we obtain that on  $\text{supp} q_{a,R}$ ,  $j_b(x)$  is equal to

$$\int J_b^a(\frac{x^a}{C} + y^a) \phi(y) dy,$$

which is a function similar to  $j_b$  if we replace  $X$  by  $X^a$ , the mollifier  $\phi$  by

$$\phi^a(x^a) = \int_{X_a} \phi(x) dx_a,$$

and (see (2.6)) the constant  $C$  by  $Cq^{1-t_a}$ . Using (2.7), this completes the proof of the Proposition.  $\square$

### 3 The resonances

In this section we describe the spectrum of the distorted Hamiltonian  $H_\theta$  and define the resonances as the discrete eigenvalues of  $H_\theta$ . We show that the resonances are the poles of the meromorphic continuation of matrix elements of the resolvent  $\langle \phi, (z - H)^{-1} \psi \rangle$  for suitable analytic vectors  $\phi, \psi$ .

As in [Hu], we denote by  $F$  the space of entire functions in  $\mathbb{C}^n$  which decay faster than any power of  $\langle z \rangle$  in some cone

$$\{z \in \mathbb{C} \mid |Imz| \leq \epsilon \langle Rez \rangle\}$$

for some  $\epsilon > 0$ . We define the set  $A$  of analytic vectors by

$$A := \{f \in L^2(X) \mid f(x) = \psi(x) \text{ for some } \psi \in F\}.$$

As in [Hu], one has :

**Lemma 3.1** *i) for any  $f \in A$ , the map*

$$\theta \mapsto U_\theta f \in L^2(X)$$

*is analytic in  $K := \{\theta \in \mathbb{C} \mid |Re\theta| \leq \epsilon_0, |Im\theta| \leq \epsilon_1\}$ .*

*ii) for any  $\theta \in K$ , the image of  $A$  under  $U_\theta$  is dense in  $L^2(X)$ .*

We first analyse the spectral properties of  $H_\theta$ .

**Theorem 3.2** *i)  $H_\theta$  with domain  $H^2(X)$  is closed and one has :*

$$\|u\|_{H^2(X)} \leq C(\|H_\theta u\| + \|u\|), \forall u \in H^2(X).$$



ii)  $H_\theta$  is  $m$ -sectorial with a sector :

$$S = \{z \in \mathbb{C} \mid |\arg(z)| \leq b < \pi/2\}.$$

iii) the essential spectrum of  $H_\theta$  is equal to

$$\sigma_{\text{ess}}(H_\theta) = \bigcup_{\alpha=2} \sigma(H_\theta^\alpha) + \frac{1}{(1+\theta)^2} \mathbb{R}^+.$$

**Definition 3.3** The points in  $\sigma(H_\theta) \setminus \sigma_{\text{ess}}(H_\theta)$  are called the resonances of  $H$ .

*proof.* let us first prove i) and ii). Since  $V_\theta$  is a bounded operator, it suffices to prove the corresponding statements with  $H_\theta$  replaced by  $H_{0,\theta} = \frac{1}{2}U_\theta D_x^2 U_\theta^{-1}$ . The Hamiltonian  $H_{0,\theta}$  is a second order differential operator with principal symbol equal to :

$$T(x, \xi) = (A(x)\xi, \xi),$$

where  $A(x) = B^{-1}(x)$ ,  $B(x) = {}^t J J$ , and  $J = 1 + \theta \nabla v$ . So  $A$  is diagonal in a basis of eigenvectors of the selfadjoint matrix  $\nabla v$  with eigenvalues :

$$(3.1) \quad 1 + 2\theta\lambda_i + 2\theta^2\lambda_i^2,$$

where  $\lambda_1 \leq \dots \leq \lambda_n$  are the eigenvalues of  $\nabla v$ . So we see that for  $|\theta| \leq \epsilon_0$  :

$$|T(x, \xi)| \geq |\xi|^2,$$

which proves i) by standard elliptic theory. To prove ii), we remark that by (3.1) the principal symbol  $T(x, \xi)$  takes its values in a convex cone strictly included in  $S$ . Then ii) follows from this observation using for example Gårding's inequality (see [Hö]).

Let us now prove iii) by induction on the number of particles. For  $N = 2$ , iii) is proven in [Hu]. Let us assume that Theorem 3.2 holds for all  $M$ -particle Hamiltonians with  $M \leq N - 1$ . For  $a \neq a_{\max}$ , let us denote by  $H_\theta^a$  the distorted Hamiltonian on  $L^2(X^a)$  obtained with the vector field  $v^a$ , and by  $\tilde{H}_{a,\theta}$  the Hamiltonian :

$$\tilde{H}_{a,\theta} := H_\theta^a + \frac{1}{2(1+\theta)^2} D_a^2.$$

Using the partition of unity defined in (2.5) and Proposition 2.3, we obtain :

$$H_\theta = \sum_{\alpha=2} q_{\alpha,R} \tilde{H}_{\alpha,\theta} + q_{a_{\max},R} H_\theta + \sum_{\alpha=2} I_{\alpha,\theta} q_{\alpha,R}.$$

Using the fact that  $H_\theta^a$  is  $m$ -sectorial by the induction hypothesis and Ichinose's lemma (see [R-S]), we get that :

$$\sigma(\tilde{H}_{a,\theta}) = \sigma(H_\theta^a) + \frac{1}{(1+\theta)^2} \mathbb{R}^+.$$

By the induction hypothesis we also have :

$$(3.2) \quad \|u\|_{H^2(X)} \leq C(\|\tilde{H}_{a,\theta}u\| + \|u\|), \quad \forall u \in H^2(X),$$

which shows as in [Hu] that :

$$\sigma_{\text{ess}}(H_\theta) = \bigcup_{|a|=2} \sigma(H_\theta^a) + \frac{1}{(1+\theta)^2} \mathbf{R}^+.$$

This completes the proof of the Theorem.  $\square$

Finally we can identify the resonances of  $H$  with poles of the meromorphic continuation of the resolvent.

**Theorem 3.4** *For any  $\phi, \psi \in A$  the quantity  $\langle \psi, (z - H)^{-1} \phi \rangle$  extends meromorphically in  $z$  from  $\{z \in \mathbf{C} \mid \text{Im} z > 0\}$  to  $\mathbf{C} \setminus \sigma_{\text{ess}}(H_\theta)$  with poles at the resonances of  $H$ . One has :*

$$\sigma_{\text{disc}}(H_\theta) = \bigcup_{\phi, \psi \in A} \{ \text{poles of } \langle \psi, (z - H)^{-1} \phi \rangle \}.$$

*If  $\lambda$  is a discrete eigenvalue of  $(H_{\theta_0})$ , and  $\theta$  varies continuously,  $\lambda$  remains a discrete eigenvalue of  $H_\theta$  as long as  $\lambda$  stays in  $\mathbf{C} \setminus \sigma_{\text{ess}}(H_\theta)$ .*

*proof.* the proof is exactly the same as in [Hu, Thm 4].  $\square$

## References

- [Ag] S. Agmon: *Lectures on exponential decay of solutions of second order elliptic equations*, Princeton University Press, Princeton 1982.
- [Ag-Co] J. Aguilar-J.M. Combes: A class of analytic perturbations for one-body Schrödinger operators, *Comm. in Math. Phys.* 22 (1971) 269-279.
- [Ba-Co] E. Balslev-J.M. Combes: Spectral properties of many body Schrödinger operators with dilation analytic potentials, *Comm. in Math. Phys.* 22 (1971) 280-294.
- [Cy] H.L. Cycon: Resonances defined by modified dilations, *Helv. Phys. Acta* 58 (1986) 969-981.
- [C.F.K.S] H.L. Cycon, R. Froese, W. Kirsch, B. Simon : *Schrödinger operators with applications to Quantum Mechanics and Global Geometry*, Texts and Monographs in Physics, Springer 1987.
- [Gr] G.M. Graf: Asymptotic completeness for  $N$ -body short range quantum systems : A new proof, *Comm. in Math. Phys.* Vol 132 (1990), 73-101.

- [Hö] L. Hörmander: *Linear partial differential operators*, Berlin Springer Verlag 1963.
- [Hu] W. Hunziker: Distortion analyticity and molecular resonances curves, *Ann. Inst. Henri Poincaré* 45 (1986), 339-358
- [R-S] M. Reed-B. Simon: *Methods of modern mathematical physics* Vol IV, New York Academic Press 1978.
- [S] I.M. Sigal: Complex transformations method and resonances in one- body quantum systems, *Ann. Inst. Henri Poincaré* 41 (1984), 103- 114.
- [Si] B. Simon: The definition of molecular resonance curves by the method of exterior complex scaling, *Phys.Lett.* 71A (1979) 211-214.