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**MAGNETIC PROPERTIES OF THE HUBBARD MODEL
WITH INFINITE INTERACTION**

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Abstract. The problem of the ground state multiplicity of the Hubbard model with the infinite interaction is considered for the systems of segments of length n .

Introduction.

The simplest model of strong correlated systems is the Hubbard model with the infinite repulsion described by the Hamiltonian

$$H = \sum_{ij} t_{ij} (c_{i\sigma}^+ c_{j\sigma} + c_{j\sigma}^+ c_{i\sigma}) (1 - c_{i,-\sigma}^+ c_{i,-\sigma}) (1 - c_{j,-\sigma}^+ c_{j,-\sigma}) \quad (1)$$

One of the most important problems is to find the dependence of the ground state multiplicity on the electron density ρ . (Nagaoka problem [1]). We consider this problem for special kind of systems, consisting of ladders of n -sites segments with different intra- and intersegment hopping integrals, t and t_{\perp} .

Spin Hamiltonian (SH).

If $\alpha = t_{\perp}/t \ll 1$ one can make use of the perturbation theory in α . For $t_{\perp} = 0$ energetic levels are spin degenerated. The resolve of degeneracy (up to α^2) leads to Hamiltonian acting upon the spin variables of the neighbouring segments [2]. It turns out that the energy of two neighbouring segments with equal (inequal) non-zero numbers of electrons is minimal at total minimal (maximal) spin. Thus, SH describes the competing interactions of ferro (F)- and antiferromagnetic (AF) types. For $\rho < 1/n$ there cannot be more than one electron at each segment and the ground state is singlet. For $\rho > 1/n$ the competition interactions leads to a separation of phases with F and AF ordering. For example, for the system with $n=2$ (ladders of dimers) and for $\rho > 1/2$ F-polaron is created. Its size is proportional to $\alpha^{-1/3}$, and at $\rho_c = (2\pi)^{-1} (3\pi\alpha \ln 2)^{1/3}$ is equal to the ladder length. Thus, the system is saturated ferromagnet. For the case with $n > 2$ polarons may be both F-and AF-type leading to the "cascade" of transitions with the alteration of total spin. Equally to the ladders a more complicated system may be considered. It

consists of dimers packed up a 2d lattice. For $\rho \leq 1/2$ SH is equivalent to Hamiltonian of the 2d t-J model which magnetic properties have been discussed in [3]. As in the case of the ladder model for $\rho > 1/2$ the F polaron appears as well as the critical density ρ_c above which the system becomes a saturated ferromagnet. The variational approach.

For ladders with arbitrary α we seek the wave function Ψ in the following form (for simplicity we consider the model with $n=2$):

$$\Psi = \prod_i N_{i\alpha} N_{i\beta} (x_i^+ y_{i\alpha} \psi_{\alpha\alpha}(i) + y_{i\beta} \psi_{\beta\beta}(i) + \psi_{\alpha\beta}(i) b_{i\alpha}^+ b_{i\beta}^+ + b_{i\alpha}^+ \psi_\beta(i) + b_{i\beta}^+ \psi_\alpha(i)),$$

where

$$\psi_{\sigma\sigma}(i) = c_{1\sigma}^+ c_{2\sigma}^+; \quad \psi_{\alpha\beta}(i) = \lambda_s \varphi_s(i) + \lambda_t \varphi_t(i); \quad \varphi_{s,t} = (2)^{-1/2} (c_{1\alpha}^+ c_{2\beta}^+ \mp c_{1\beta}^+ c_{2\alpha}^+); \quad \psi_\sigma(i) = (2)^{-1/2} (c_{1\sigma}^+ + c_{2\sigma}^+).$$

1,2 are the numbers of sites of a i -th segment; λ_s, λ_t are the variational parameters at $\lambda_s^2 + \lambda_t^2 = 1$; $x_i^2 + y_{i\alpha}^2 + y_{i\beta}^2 = 1$ ($y_{i\sigma} = 0$ or 1). $N_{i\sigma} = b_{i\sigma}^+ b_{i\sigma}$; $b_{i\sigma}^+$ are the operators of pseudo-fermi-particles.

Effective Hamiltonian for pseudoparticles is $H_{\text{eff}} = \langle \Psi | H | \Psi \rangle$ where $\langle \dots \rangle$ means an average on c -operators. Analysing H_{eff} we have reproduced the results obtained by SH method for $\alpha \ll 1$. In isotropic case, $\alpha=1$, and for $\rho < 1/2$ a ground state is singlet. At $\rho \rightarrow 1$ a ground state spin is maximal. At $\rho=0.67$ this state becomes unstable with respect to overturning single up-spin electron. However, it turns out that at $\rho=0.73$ F state is unstable against multimagnon excitations. Note that in Gutzwiller approximation F state in this model is stable at all values of ρ .

We also studied the dependence $S(\rho)$ using the exact diagonalization of small ladders of dimers (up to 16 sites). The agreement between analytical and numerical results is fairly good.
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