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# Neutron scattering investigations of the $\text{LaNi}_5$ hydrogen storage system

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*Abstract.* Static and dynamic properties of hydrogen in  $\text{LaNi}_5\text{H}_x$  ( $0 \leq x \leq 6$ ) compounds were determined by neutron scattering. Deuterium positions as well as distribution have been determined by neutron diffraction. The neutron inelastic scattering technique has been used to study the hydrogen diffusion process from the dependence of the quasielastic line width upon momentum transfer. The results are in qualitative agreement with the predictions of a jump diffusion model, and the diffusion parameters are qualitatively discussed in terms of the structure information.

## 1. Introduction

Metal hydrides and in particular  $\text{LaNi}_5\text{H}_x$  [1] are important with respect to technical applications such as use for hydrogen storage, thermal compression, in electrochemical cells, for catalysis or hydrogen purification.

By means of neutron diffraction investigations hydrogen may be located on an atomic scale (crystal structure) and inelastic neutron scattering experiments provide primary information about the dynamics of hydrogen and in particular concerning hydrogen diffusion. Thus neutron scattering contributes essentially to the understanding of physical properties of metal hydrides.

Preliminary neutron diffraction results are published for  $\text{LaNi}_5\text{D}_6$  [2] and  $\text{LaNi}_5\text{D}_7$  [3]. In the former study only hexagonal lattice constants (cf. Table 1) were derived, although the presence of pseudo-hexagonal twinning due to an orthorhombic structure similar to  $\text{PrCo}_5\text{D}_x$  is assumed. Bowman et al. [3] proposed a structure based on the trigonal space group  $P31m$ , which corresponds to two sets of hydrogen sites. Thus hydrogen storage is associated with symmetry reduction compared to the  $P6/mmm$  space group symmetry of  $\text{LaNi}_5$ . Because of experimental difficulties such as dominant peaks due to the steel sample container the derived structure appears to be rather uncertain.

Various structure models based on space group  $P6/mmm$  were tested in a proton nuclear magnetic resonance (NMR) study of  $\text{LaNi}_5\text{H}_6$  [4], where moreover hydrogen diffusion (cf. also [5]) was investigated. It is concluded that a  $\text{PrCo}_5\text{H}_4$  like structure has to be rejected for  $\text{LaNi}_5\text{H}_6$ . Later NMR measurements of  $\text{LaNi}_5\text{H}_x$  and  $\text{LaNi}_5\text{D}_x$  [6] confirm the structure proposed by Bowman et al. [3].

In order to obtain more precise information concerning hydrogen distribution and diffusion neutron scattering experiments were made on deuterated or hydrogenated powder samples of  $\text{LaNi}_5$ , which do not show magnetic order (Pauli paramagnetism) [2,7]. In the subsequent sections of this paper we describe after discussion of experimental details the structures. In the last part measurements concerning diffusion are presented and discussed with respect to structure. (Concerning preliminary results cf. [8].)

## 2. Experimental details

$\text{LaNi}_5$  was prepared from 99.9% pure La (Research Chemicals) and 99.998% pure Ni (Koch Light) by induction melting in a water cooled copper boat in high vacuum ( $2 \times 10^{-6}$  Torr) and subsequent quenching. X-ray analysis yields hexagonal  $\text{CaCu}_5$  type structure for  $\text{LaNi}_5$  (concerning lattice constants cf. Table 1).

Pauli paramagnetism of  $\text{LaNi}_5$  was verified by measurements of magnetic susceptibility. In order to obtain samples with temperature ( $70 \div 300$  K) and field ( $0 \div 4$  kOe) independent magnetic susceptibility it was necessary to compensate for small evaporation losses and possible impurity effects of the starting materials by reducing the initial Ni content corresponding to  $\text{LaNi}_{4.97}$  (susceptibility  $4.6 \times 10^{-6}$  emu/g).

The samples were hydrogenated or deuterated in stainless steel containers at pressure of 50 bars. Hydrogenation took place within a few minutes and yielded the composition  $\text{LaNi}_5\text{H}_6$  ( $\text{LaNi}_5\text{D}_6$ ).

For the elastic (inelastic) neutron scattering experiments the samples were enclosed in cylindrical vanadium (aluminium) containers of 1 cm diameter at room temperature under constant deuterium (hydrogen) pressure up to 5 bars. The deuterium (hydrogen) content of  $\text{LaNi}_5$  samples was measured by desorption after the neutron experiments. As hydrogenation generally results in very fine powders,  $\text{LaNi}_5$  was measured after hydrogenation and subsequent desorption.

The neutron diffraction experiments were performed on two-axis spectrometers at reactor Diorit, Würenlingen with neutrons of wavelength  $\lambda = 2.312$  and  $1.192$  Å, using graphite 002 and Ge 311 as monochromators respectively. The observed intensities were corrected for absorption and incoherent scattering according to the measured transmission of the samples.

The inelastic neutron experiments were carried out on a triple-axis spectrometer at the reactor Diorit. The incoming neutrons filtered by cooled beryllium had an energy of 5 meV. Since the observed line widths due to hydrogen diffusion were rather small, the experiments have been repeated on a MARC spectrometer with improved energy resolution (incoming neutron energy: 4.6 meV) and larger intensity. The agreement between the resulting line widths obtained from the two measurements was satisfactory.

## 3. Localization of Deuterium

The final diffraction patterns were analyzed by means of the profile method [9], using the neutron scattering lengths  $b_{\text{La}} = 8.3$  F,  $b_{\text{Ni}} = 10.3$  F and  $b_{\text{D}} = 6.67$  F.

Neutron diffraction measurements were made at room temperature on powder samples of  $\text{LaNi}_5$ ,  $\text{LaNi}_5\text{D}_{0.15}$  ( $\alpha$ -phase) and  $\text{LaNi}_5\text{D}_6$  ( $\beta$ -phase). The composition  $x$  corresponds to desorption measurements. Results are summarized in Tables 1 to 3.

Table 1

Structure parameters of  $\text{LaNi}_5\text{D}_x$ . Standard deviations are given within parentheses and refer to the last digit.  $N$  denotes neutron diffraction. For  $x = 0, 0.2$ : space group  $P6_3/\text{mmc}$  (atom positions: La: 1a, Ni1: 2c, Ni2: 3g, D: 3f); for  $x = 6$ : space group  $P3_1\text{m}$  (atom positions: La: 1a, Ni1: 2b, Ni2: 3c, D1: 3c, D2: 6d).  $B$  = Debye-Waller parameter (intensity  $\sim e^{-2W}$ ,  $W = B(\sin \theta/\lambda)^2$ ).  $R$ ,  $RW_p$  = agreement values concerning integrated and profile intensities, respectively [9].

$x$	0	0.2	6						$\gamma$ [3]
	(a)	(b)	(c)	(a)	(b)	(c)	(d)	(e) [2]	
$\lambda$ [Å]	2.312	2.312	2.312	2.312	2.312	2.312	2.312	2.57	1.142
$a_N$ [Å]	5.017(5)	5.025(5)	5.025(5)	5.410(5)	5.410(5)	5.410(5)	5.410(5)	5.393	5.387
$c_N$ [Å]	3.986(5)	3.988(5)	3.988(5)	4.293(5)	4.293(5)	4.293(5)	4.293(5)	4.285	4.273
positional parameters									
$z_{\text{La}}$	0	0	0	0	0	0	0		0
$z_{\text{Ni1}}$	1	1	1	0.94(1)	0.93(1)	0.93(1)	0.94(1)		0.95
$x_{\text{Ni2}}$	0.5	0.5	0.5	0.480(2)	0.480(2)	0.481(3)	0.476(3)		0.50
$z_{\text{Ni2}}$	0.5	0.5	0.5	0.482(8)	0.479(8)	0.481(7)	0.483(7)		0.48
$x_{\text{D1}}$	0.5	0.5	0.5	0.470(3)	0.471(3)	0.469(3)	0.477(5)		0.50
$z_{\text{D1}}$	0	0	0	0.077(7)	0.076(7)	0.076(8)	0.081(7)		0.08
$x_{\text{D2}}$				0.180(4)	0.176(4)	0.179(4)	0.182(4)		0.25
$y_{\text{D2}}$				0.832(4)	0.828(4)	0.831(4)	0.832(4)		0.86
$z_{\text{D2}}$				0.56(2)	0.55(2)	0.56(2)	0.56(1)		0.58
deuterium distribution									
D1	0.26(2)	0.15*	0.24(2)	3.0(1)	2.86(4)	2.84(4)	3.6(6)		3.0
D2				3.4(1)	3.14	3.16	4.2(6)		3.8
$B(B_{\text{La, Ni}})$ [Å <sup>2</sup> ]	1.1(1)	1.3(1)	0.86(4)	2.0(2)	2.0(2)	2.2(3)	1.0(7)		
$B_D$						1.6(3)	4(1)		
$R$	0.037	0.053	0.067	0.054	0.054	0.054	0.054		
$RW_p$	0.145	0.135	0.179	0.173	0.173	0.173	0.172		

\*) Values corresponding to desorption measurements.



Table 3  
Observed, absorption corrected and calculated integrated intensities of  $\text{LaNi}_5\text{D}_6$  (cf. a, Table 1).

$h\ k\ l$	$I_{\text{obs}}$	$\pm$	$\Delta I_{\text{obs}}$	$I_{\text{calc}}$	$h\ k\ l$	$I_{\text{obs}}$	$\pm$	$\Delta I_{\text{obs}}$	$I_{\text{calc}}$
1 0 0	53178		347	51064	2 1 1	149		12	115
0 0 1	3065		250	2948	2 1 $\bar{1}$	2208		181	1705
1 0 1	0		256	292	2 0 2	50868		422	52740
1 1 0	10764		294	9174	3 0 0	1679		206	1336
2 0 0	64608		339	63173	3 0 1	64168		492	62918
1 1 1	136887		446	129343	0 0 3	2278		262	3559
0 0 2	84921		435	87889	1 0 3	6298		64	5896
2 0 1	32242		354	31736	2 1 2	916		7	863
1 0 2	5580		344	9596	2 1 $\bar{2}$	13932		104	13127
2 1 0	18244		392	12446	2 2 0	176757		573	173186
1 1 2	8617		337	10451	3 1 0	21770		411	24070
					2 2 1	1995		50	2174

Illustrative diffraction patterns are shown in Fig. 1. Apparently hydrogen storage in  $\text{LaNi}_5$  is associated with essential, reversible modifications of the hexagonal  $\text{CaCu}_5$  type structure of  $\text{LaNi}_5$  (cf. also Tables 2 and 3), yielding ternary compounds. Considerable line broadening due to small grain size of the  $\text{LaNi}_5\text{D}_x$  powders as a consequence of hydrogen absorption and desorption raised difficulties in least squares fitting of the profile intensities. (Despite similar instrumental conditions the peak halfwidths of e.g.  $\text{La}_7\text{Ni}_3$  were found to be roughly 50% smaller.)

Positions and intensities of the measured\*  $\text{LaNi}_5$  sample confirm the lattice constants ( $a_x = 5.01^*$ ,  $5.017\text{ \AA}$  [1];  $c_x = 3.98^*$ ,  $3.987\text{ \AA}$  [1]) and the  $\text{CaCu}_5$ -type structure (cf. Table 1).

At low deuterium concentrations ( $\alpha$ -phase) the lattice expands slightly in the basal plane ( $\Delta a/a \sim 0.16\%$ ), whereas  $c$  remains almost constant ( $\Delta c/c \sim 0.05\%$ ). Similar X-ray results are mentioned in [2]. In the  $\beta$ -phase ( $\text{LaNi}_5\text{D}_6$ ) the unit cell is expanded 8% and 25% with respect to lattice constants and volume respectively – almost without change of the  $c/a$  ratio – compared to  $\text{LaNi}_5$ .

In the  $\alpha$ -phase ( $\text{LaNi}_5\text{D}_{0.25 \pm 0.02}$  according to neutron diffraction, in reasonable agreement with the composition  $x = 0.15$  determined by desorption, cf. Table 1) best agreement of observed and calculated intensities was obtained for space group  $\text{P6}/\text{mmm}$  ( $\text{LaNi}_5$ ), assuming a statistical distribution of deuterium on or at least “centered” on positions  $3f$  ( $\frac{1}{2} 0 0$ ,  $0 \frac{1}{2} 0$ ,  $\frac{1}{2} \frac{1}{2} 0$ ). The small deuterium concentration renders a precise localization of deuterium difficult, because the contribution to the structure factor is correspondingly small ( $R = 0.069$  at  $\lambda = 2.312\text{ \AA}$ , neglecting deuterium). Therefore it is intended to perform also measurements on a  $\text{LaNi}_5\text{D}_x$  sample with maximum deuterium content in the  $\alpha$ -phase ( $x \sim 0.5$  [1]). Considering the structure model, a main lattice expansion within the basal plane is expected. Presumably the local distortions around deuterium atoms are larger than the average distortions reflected in the lattice constants.

Concerning the  $\beta$ -phase ( $x = 6$ ) the interpretation of the measurements does not appear to be possible within the context of the  $\text{P6}/\text{mmm}$  space group of  $\text{LaNi}_5$ , in agreement with previous investigations [3, 6]. Reasonable fits of the measured integrated intensities were obtained for space group  $\text{P31m}$  ( $R \sim 0.05$ , cf. Table 1) and less satisfactory agreement results for space groups  $\text{P3}$  (deuterium D on sites  $3d$  and  $3d'$ ,



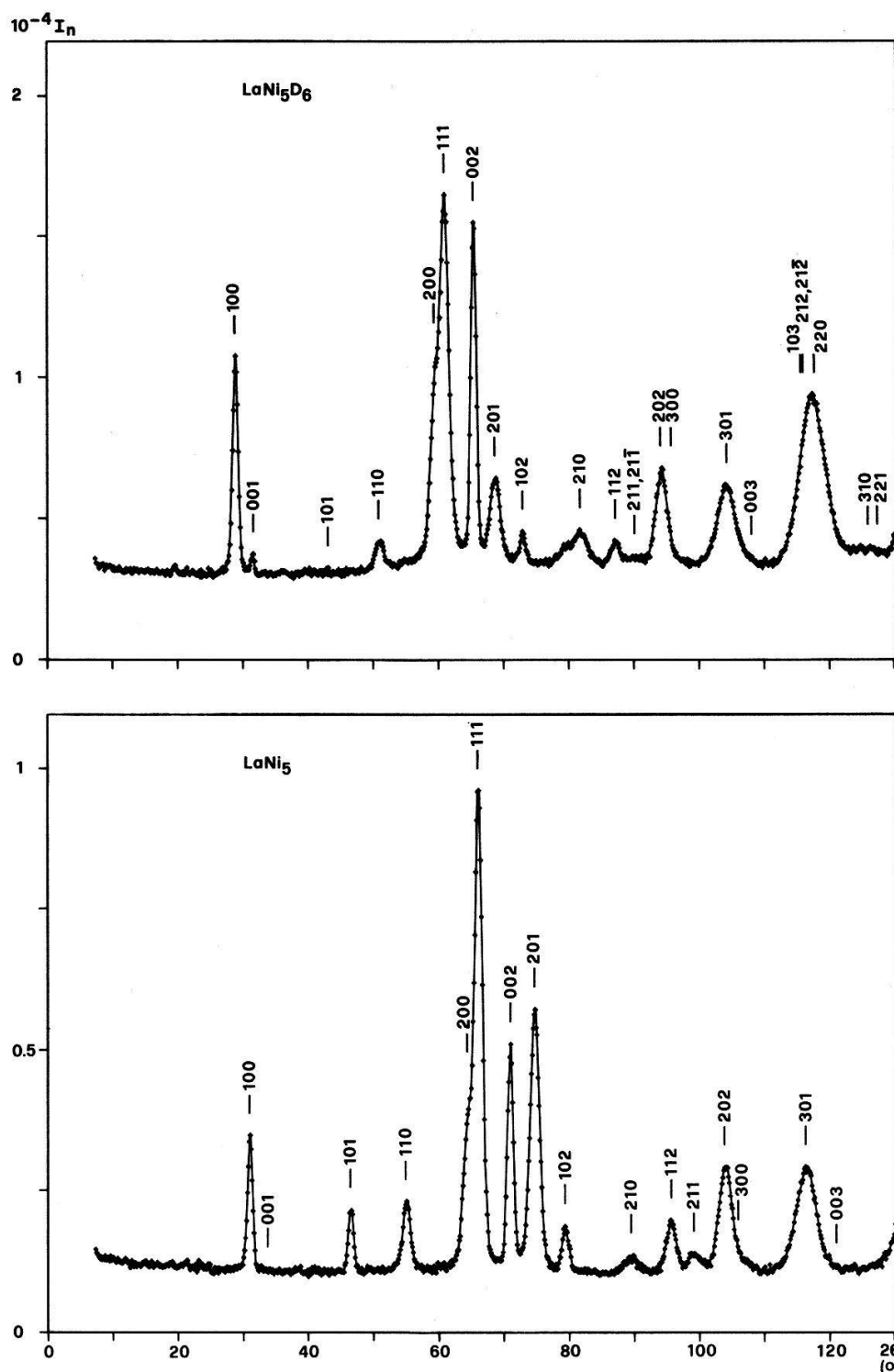


Figure 1  
Neutron diffraction patterns (absorption corrected) of  $\text{LaNi}_5$  and  $\text{LaNi}_5\text{D}_6$ .  $2\theta$  = scattering angle,  $I_n$  = neutron intensity.

$R \sim 0.08$ ) and  $P3m1$  (D on 3d and 3d',  $R \sim 0.10$ ). Results of calculations based on different numbers of parameters are compared in Table 1 (a ÷ d), which yield the same  $R$  value. The neutron intensities indicate a deuterium content  $x = 6.4 \pm 0.2$ , in reasonable agreement with the value 6.0 determined by desorption. Positional parameters and deuterium distribution do not depend essentially on the assumptions made concerning thermal motion, except for d, Table 1, where the too large  $x$  value is

caused by correlations with the temperature factor of deuterium. Insertion of large amounts of deuterium in  $\text{LaNi}_5$  leads to a marked distortion of the Ni configuration (cf. Table 1). The crystal structure (space group  $P31m$ ) of  $\text{LaNi}_5\text{D}_6$  is shown in Fig. 2. It corresponds to the model proposed by Bowman et al. [3], which implies a statistical distribution of approximately 3 deuterium atoms on tetrahedral sites 6d (D2) and completely occupied (3D) tetrahedral sites 3c (D1). The corresponding smallest Ni–D distances of 1.45 to 1.74 Å are remarkably short. Of reasonable magnitude appear the smallest La–D distances of 2.50 to 2.91 Å. Concerning hydrogen diffusion the shortest D–D distances of 1.57 and 1.69 Å (D2–D2), 2.48 Å (D1–D2), 2.59 Å (D2–D2), 2.61 Å and 2.68 Å (D1–D2), 2.72 Å (D1–D1), 2.79 Å (D1–D2) as well as 2.83 Å and 3.06 Å (D2–D2) are presumably of importance. Seemingly sites 3c (corresponding approximately to positions 3f with respect to space group  $P6/mmm$ ) are first filled during hydrogen storage. Finally interstices 3c are 100% and positions 6d 57% occupied in  $\text{LaNi}_5\text{D}_6$ .

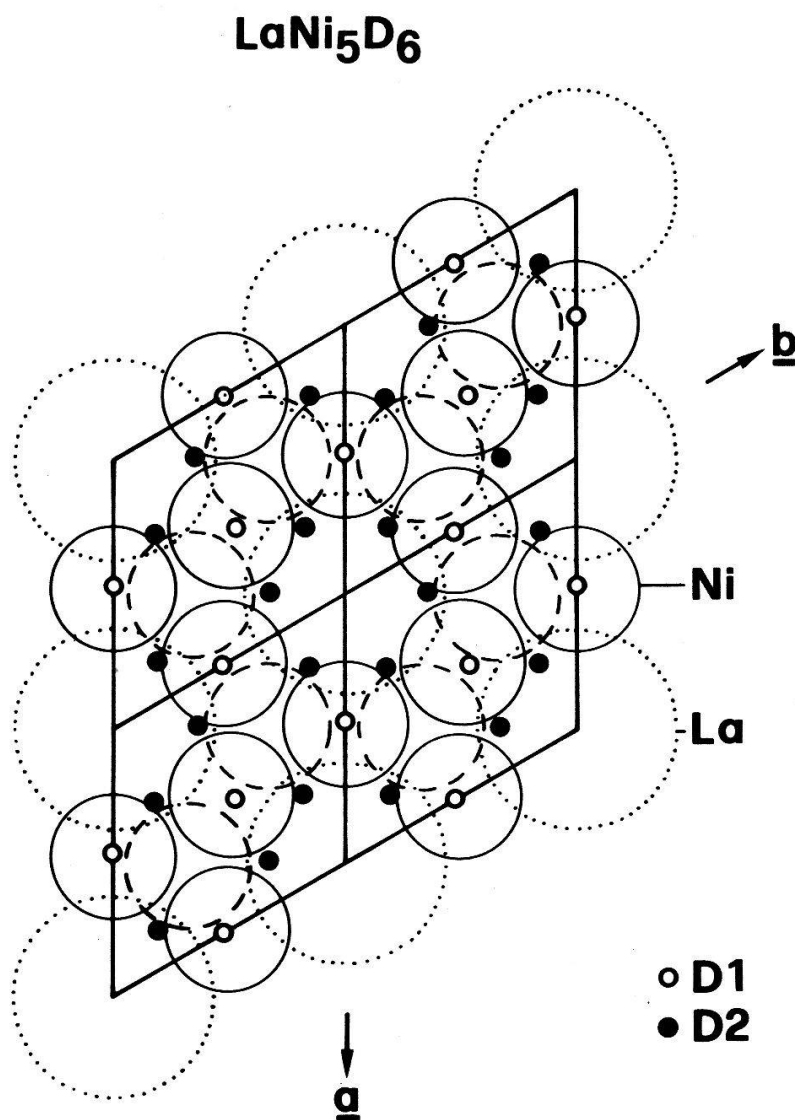


Figure 2

Projection of the crystal structure of  $\text{LaNi}_5\text{D}_6$  (space group  $P31m$ ) on the basal plane perpendicular to  $[001]$  (basic translations  $a$ ,  $b$ ). Four units cells are shown. Atomic radii of La ( $\cdots$ ;  $z = 0$ ) and Ni ( $---$ ;  $z = -0.06$ ,  $---$ ;  $z = 0.482$ ) are estimated on the basis of the  $\text{LaNi}_5$  structure ( $r_{\text{La}} \sim 1.99$  Å,  $r_{\text{Ni}} \sim 1.25$  Å). Open ( $z = 0.077$ ) and filled circles ( $z = 0.56$ ) indicate approximately 100 and 50% occupied deuterium positions respectively.



#### 4. Hydrogen diffusion in $\text{LaNi}_5\text{H}_x$

The neutron inelastic scattering technique has been used to study the dynamics of hydrogen in  $\text{LaNi}_5\text{H}_x$  ( $x = 0, 0.2, 6$ ) at room temperature. Such measurements can provide information about the diffusion constant, diffusion rates and activation energies. Furthermore, in some cases hydrogen jump distances and directions and thus site occupations can be determined from the dependence of the quasielastic line width  $\Gamma$  upon momentum transfer  $\hbar Q$ .

In the data analysis the scattering contributions of  $\text{LaNi}_5$  have been subtracted from the observed energy spectra of  $\text{LaNi}_5\text{H}_x$  by taking account of the different transmission factors due to hydrogenation. Then the resulting quasielastic peaks have been fitted to a single Lorentzian folded over the incident neutron spectrum, with the width and height as fitting parameters. The measured widths  $\Gamma$  for  $x = 6$  are shown in Fig. 3 which displays a broad maximum at  $Q = 1.6 \text{ \AA}^{-1}$ . The experiments for  $x = 0.2$  are not yet completed, but the results obtained so far indicate a similar maximum shifted to  $Q = 2.3 \text{ \AA}^{-1}$ .

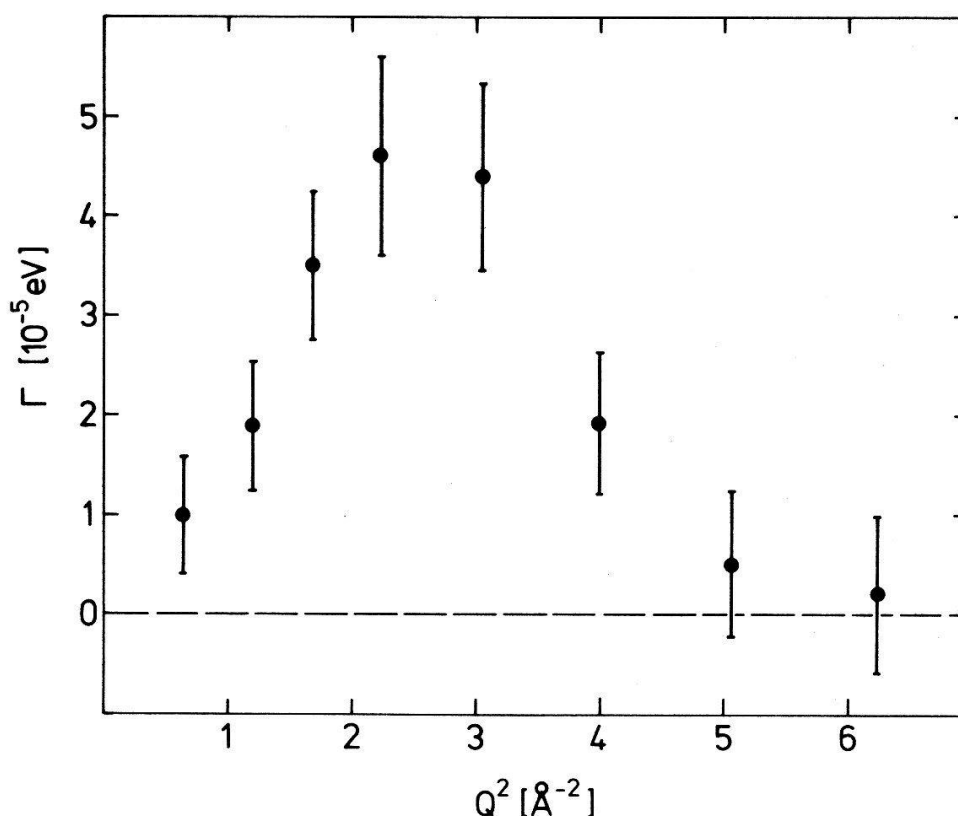


Figure 3

Widths  $\Gamma$  of the quasielastic line for neutron scattering from hydrogen in  $\text{LaNi}_5\text{H}_6$  as a function of  $Q^2$  ( $\hbar Q$  = momentum transfer) at room temperature.

Our results are in qualitative agreement with the predictions of a jump diffusion model, supposing that the diffusion process consists of thermally activated steps which are statistically independent. This assumption is likely to be justified if the hydrogen atom performs many oscillations between successive steps. Furthermore, we neglect the time needed for a jump from site to site compared to the mean time of stay  $\tau$  ( $1/\tau$  = mean jump rate). Finally we assume that correlation effects do not affect the jump probability. This may hold even for the high hydrogen concentration

in  $\text{LaNi}_5\text{H}_6$ , since there are many unoccupied interstitial positions available (see Fig. 2). The jump diffusion model yields [10]

$$\Gamma = \frac{2\hbar}{n\tau} \sum_{j=1}^n (1 - \exp \{-i\mathbf{Q} \cdot \mathbf{R}_j\}), \quad (1)$$

where  $\mathbf{R}_j$  is the jump vector. From the maximum in Fig. 3 we obtain for  $\text{LaNi}_5\text{H}_6$  jump lengths of the order of 4 Å which are considerably larger than the shortest H–H distances. On the other hand the preliminary results for  $\text{LaNi}_5\text{H}_{0.2}$  suggest jump distances of  $2.5 \div 3$  Å in agreement with the shortest H–H distances given in Section 3.

For  $Q \rightarrow 0$  equation (1) approaches

$$\Gamma = 2\hbar D Q^2, \quad (2)$$

i.e. the diffusion constant  $D$  can be determined from the initial slope of  $\Gamma$  vs.  $Q^2$ . From Fig. 3 we obtain for  $\text{LaNi}_5\text{H}_6$   $D = (1.2 \pm 0.5) \times 10^{-6} \text{ cm}^2\text{s}^{-1}$  which is an order of magnitude larger than the value obtained from NMR studies [4]. Results for  $\text{LaNi}_5\text{H}_{0.2}$  are not yet available in the low  $Q$  region.

## 5. Conclusions

The main features of deuterium storage in hexagonal  $\text{LaNi}_5\text{D}_x$  compounds have been established with respect to crystal structure by means of neutron diffraction. Associated with the formation of ternary compounds the lattice expands, and reversible modifications of the  $\text{CaCu}_5$ -type structure of  $\text{LaNi}_5$  occur. In the  $\alpha$ -phase ( $x \sim 0.2$ ) deuterium is distributed on sites 3f (space group  $\text{P6}/\text{mmm}$  as for  $x = 0$ ), corresponding to primary expansion in the basal plane. In the case of large deuterium concentration ( $\beta$ -phase,  $x \sim 6$ ) the unit cell is expanded 25% with respect to volume, almost with the same  $c/a$  ratio as for  $\text{LaNi}_5$ . The structure parameters of the model based on space group  $\text{P31m}$  [3] are refined. The Ni configuration is considerably distorted compared to  $\text{LaNi}_5$ . Deuterium fills tetrahedral interstices 3c and occupies 57% of tetrahedral sites 6d according to a statistical distribution.

The inelastic neutron scattering results provide information concerning the dynamics of hydrogen in the hydrogen storage system  $\text{LaNi}_5\text{H}_x$ . The data are qualitatively discussed in terms of a jump diffusion model and related to the structure. The diffusion constant  $D = (1.2 \pm 0.5) \times 10^{-6} \text{ cm}^2\text{s}^{-1}$  is an order of magnitude larger than the value obtained from NMR studies [4]. From the size of the diffusion constant we conclude that diffusion does not control the hydrogen sorption rates in  $\text{LaNi}_5$  compounds.

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