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# Density of states in spectral geometry

TOSHIAKI ADACHI and TOSHIKAZU SUNADA

### Introduction

Recent studies on spectal geometry threw a light on the relationships between a discontinuous action of a group and the spectrum of the Laplacian (or more generally the spectra of elliptic operators) on a non-compact Riemannian manifold. The first result in this direction is the observation by R. Brooks [B] that the bottom of the  $L^2$ -spectrum of the Laplacian is related to the amenability of discontinuous transformation groups (see also [KOS] and [S1]). The purpose of this paper is to investigate the integrated density of states of a periodic Schrödinger operator on a manifold with compact quotient from the same point of view.

The integrated density of states, which is the concept introduced first by physicists in quantum theory of solids, is a non-decreasing function  $\varphi(\lambda)$  on the real line defined roughly as the number of possible energy levels in the interval  $(-\infty, \lambda)$  divided by the volume of a sufficiently large domain. To justify this physical definition, we must impose a suitable boundary condition on eigenfunctions and specify the way how to blow up the domain filling the whole space. For the Schrödinger operator with a periodic potential on the Euclidean space, a classical observation (cf. [Sh]) says that  $\varphi(\lambda)$  is well-defined as far as the domain blows up in a sufficiently regular way and does not depend on the choice of the boundary conditions; say, Dirichlet, Neumann, and periodic boundary conditions. It is also a fact that the spectrum of the Schrödinger operator on the whole space is characterized as the set of increasing points of  $\varphi(\lambda)$ . One of the results in this paper gives a partial generalization of those facts to the case of a Riemannian manifold with compact quotient. In the discussion, we shall see a prominent role of amenability of discontinuous groups acting on manifolds, together with the role of spectral distribution functions defined by means of the concept of the von Neumann trace. See [SN], [S3], and [KOS] for the general background for the spectral theory of periodic Schrödinger operators on a manifold, also [ES] which gives us the stimulus for writing this paper.

### §1. Definitions and results

Let X be a complete, connected, noncompact Riemannian manifold of dimension n, and  $\mathcal{D} = \{D_j\}_{j=1}^{\infty}$  be a family of bounded connected open sets in X with piecewise smooth boundaries satisfying

$$\bar{D}_j \subset D_{j+1}, \qquad \bigcup_{j=1}^{\infty} D_j = X.$$

Let q be a smooth real-valued function on X. Consider the Schrödinger operator  $H_{D_j} = -\Delta_{D_j} + q$  on each  $D_j$  acting on  $L^2(D_j)$  with Dirichlet boundary conditions. We denote by  $\varphi_{D_j}(\lambda)$  the number of eigenvalues of  $H_{D_j}$  not exceeding  $\lambda$ , where each eigenvalue is repeated according to its multiplicity. We now define the function  $\varphi_{\mathscr{D}}$  by the limit (if it exists)

$$\varphi_{\mathscr{D}}(\lambda) = \lim_{j \to \infty} \operatorname{vol}(D_j)^{-1} \varphi_{D_j}(\lambda),$$

and call  $\varphi_{\mathscr{D}}$  the integrand density of states for the Schrödinger operator  $H_X = -\Delta_X + q$  associated with the family  $\mathscr{D}$ . The questions with which we are concerned are: (1) Under what condition does the limit exist? (2) When it exists, is  $\varphi_{\mathscr{D}}$  independent of the choice of the expanding family  $\mathscr{D}$ ?

Given a manifold with *compact quotient*, we may introduce the integrated density of states associated with *periodic boudary conditions*. Here a complet Riemannian manifold X is said to have compact quotient if there is a discrete subgroup  $\Gamma$  in the isometry group of X acting discontinuously on X such that the quotient space  $M = \Gamma \setminus X$  is compact. We assume that q is  $\Gamma$ -invariant so that q may be regarded as a function on M. Let

$$H_X = \int \lambda \ dE(\lambda)$$

denote the spectral resolution of  $H_X$ . We define the spectral distribution function  $\Phi_{\Gamma}$  by

$$\Phi_{\Gamma}(\lambda) = \operatorname{Tr}_{\Gamma} E(\lambda),$$

where  $\operatorname{Tr}_{\Gamma}$  is the standard von Neumann trace on the von Neumann algebra of  $\Gamma$ -equivariant bounded operators of  $L^2(X)$  (see [At], [ES], [S3]).

From the definition, it follows easily that the quantity vol  $(\Gamma \setminus X)^{-1}\Phi_{\Gamma}(\lambda)$  depends only on the commensurability class of  $\Gamma$ ; that is, if  $\Gamma_1$  and  $\Gamma_2$  are

discontinuous transformation groups of X such that q is invariant under  $\Gamma_1$  and  $\Gamma_2$ , and  $\Gamma_1 \cap \Gamma_2$  is of finite index in both  $\Gamma_1$  and  $\Gamma_2$ , then one has

$$\operatorname{vol}(\Gamma_1 \backslash X)^{-1} \Phi_{\Gamma_1} = \operatorname{vol}(\Gamma_2 \backslash X)^{-1} \Phi_{\Gamma_2}.$$

In the special case that  $q \equiv 0$  and X is a homogeneous Riemannian manifold, the quantity vol  $(\Gamma \setminus X)^{-1}\Phi_{\Gamma}(\lambda)$  does not depend on  $\Gamma$ . For example, if  $X = \mathbb{R}^2$ , one has

$$\operatorname{vol}(\Gamma \setminus X)^{-1} \Phi_{\Gamma}(\lambda) = (4\pi)^{-1} \lambda, \qquad \lambda \geq 0,$$

and if  $X = \mathbb{H}^2$ , the hyperbolic 2-plane, one has

$$\operatorname{vol}(\Gamma \backslash X)^{-1} \Phi_{\Gamma}(\lambda) = (4\pi)^{-1} \int_{0}^{\lambda - 1/4} \tanh \pi \sqrt{\lambda} \, d\lambda, \qquad \lambda \geq 1/4.$$

To see that the function  $\Phi_{\Gamma}$  may be regarded as the integrated density of states associated with periodic boundary value conditions, we suppose that  $\Gamma$  acts freely on X and has a family of normal subgroup  $\{\Gamma_i\}_{i=1}^{\infty}$  such that  $\Gamma_i$  is of finite index in  $\Gamma$ ,  $\Gamma_{i+1}$  is contained in  $\Gamma_i$ , and  $\bigcap_{i=1}^{\infty} \Gamma_i = \{1\}$ . We then have a tower of finite-fold covering maps of closed manifolds

$$\cdots \longrightarrow M_{i+1} \longrightarrow M_i = \Gamma_i \backslash X \longrightarrow \cdots \longrightarrow M_1 \longrightarrow M.$$

Let  $\Phi_{M_i}(\lambda)$  denote the number of eigenvalues of  $H_{M_i}$  on the closed manifold  $M_i$  not exceeding  $\lambda$ . In [SN], it was observed that

$$\operatorname{vol}(\Gamma \backslash X)^{-1} \Phi_{\Gamma}(\lambda) = \lim_{i \to \infty} \operatorname{vol}(M_i)^{-1} \Phi_{M_i}(\lambda)$$

at all points of continuity of  $\Phi_{\Gamma}$ . It should be noted ([ES], [SN]) that the set of increasing points of  $\Phi_{\Gamma}$  coincides with the spectrum of  $H_{\chi}$ .

It is natural to compare  $\varphi_{\mathscr{D}}(\lambda)$  with  $\Phi_{\Gamma}(\lambda)$ . In the case that X is the Euclidean space  $\mathbb{R}^n$  and  $\mathscr{D}$  is a family of concentric balls, it is known that  $\varphi_{\mathscr{D}}$  exists and coincides with vol  $(\Gamma \setminus X)^{-1}\Phi_{\Gamma}$ . On the other hand, if  $X = \mathbb{H}^n$ , the n-dimensional hyperbolic space, and  $\mathscr{D}$  is a family of concentric geodesic balls in  $\mathbb{H}^n$ , we observe that  $\varphi_{\mathscr{D}}$  is not equal to vol  $(\Gamma \setminus X)^{-1}\Phi_{\Gamma}$  (see Section 3). This is due to different geometric features of geodesic balls in  $\mathbb{R}^n$  and  $\mathbb{H}^n$  which may be clarified if one looks at the ratio

$$\operatorname{vol}(\partial_h D_i)/\operatorname{vol}(D_i), \quad h > 0,$$

where  $\partial_h D$  denotes the "thick" boundary  $\{x \in D; \operatorname{dist}(x, \partial D) \leq h\}$ . In fact, for  $\mathbb{R}^n$ , this goes to zero as  $j \to \infty$  for every h, while, for  $\mathbb{H}^n$ , this goes to the positive number  $1 - e^{-h(n-1)}$ . In terms of discrete transformation groups, this corresponds to the fact that a group  $\Gamma$  acting discontinuously on  $\mathbb{R}^n$  is amenable, and a group  $\Gamma$  acting on  $\mathbb{H}^n$  is non-amenable. Indeed, we may prove the following general criterion of amenability.

**PROPOSITION** 1.1. The transformation group  $\Gamma$  is amenable if and only if there exists an expanding family  $\mathcal{D} = \{D_j\}$  of bounded domains with piecewise smooth boundaries satisfying the following property:

$$\lim_{j \to \infty} \operatorname{vol}\left(\partial_h D_j\right) / \operatorname{vol}\left(D_j\right) = 0 \tag{P}$$

for every h > 0.

In light of this criterion, we now state the main theorem of this paper, a generalization of the classical result for  $X = \mathbb{R}^n$ .

THEOREM 1.1. If an expanding family  $\mathcal{D} = \{D_j\}$  satisfies the property (P) in the above proposition, then  $\varphi_{\mathcal{D}}$  exists and equals vol  $(\Gamma \setminus X)^{-1}\Phi_{\Gamma}$  at all points of continuity of  $\Phi_{\Gamma}$ .

An immediate consequence of this theorem is that, if  $\Gamma$  is amenable, the integrated density of states  $\varphi_{\mathscr{D}}$  does not depend on the expanding family  $\mathscr{D}$  with the property (P). We also conclude that vol  $(\Gamma \setminus X)^{-1}\Phi_{\Gamma}$  does not depend on  $\Gamma$ , which is by no means trivial from the definition of  $\Phi_{\Gamma}$  since X is not supposed to be homogeneous.

It is interesting to consider the density of states associated with *Neumann boundary conditions*. We conjecture that the same statements as in Theorem 1 hold. Sznitman [Sz2] shows that, for the hyperbolic space, the integrated density of states associated wity Dirichlet boundary conditions is different from that associated with Neumann boundary conditions.

## §2. Families of expanding domains and limit relations for the heat kernels

Henceforth we assume that X is a Riemannian manifold with compact quotient  $\Gamma \setminus X$ . We choose a fundamental domain  $\mathscr{F}$  for the action of  $\Gamma$  with compact closure. The distance function on X will be denoted by d(x, y).

Let k(t, x, y) denote the heat kernel function for the semi-group  $\exp(-tH_X)$ , and  $k_D(t, x, y)$  the heat kernel function on a domain D associated with Dirichlet boundary conditions. We readily get

$$\int e^{-\lambda t} d\varphi_D(\lambda) = \int_D k_D(t, x, x) dx.$$

The following lemma on the spectral distribution function  $\Phi_{\Gamma}$  is immediate from the definition of  $\Gamma$ -trace.

LEMMA 2.1. 
$$\int e^{-\lambda t} d\Phi_{\Gamma}(\lambda) = \int_{\mathscr{F}} k(t, x, x) dx.$$

The idea of proof of Theorem 1.1 is based on a uniform estimate of the difference between the diagonal of the heat kernel and that of the Dirichlet heat kernel.

LEMMA 2.2. Given a positive T, we have positive constants  $C_1$  and  $C_2$  such that

$$0 \le k(t, x, y) \le C_1 t^{-n/2} \exp\left(-C_2 d(x, y)^2 / t\right) \tag{1}$$

for  $t \in (0, T]$ , and

$$0 \le k(t, x, y) - k_D(t, x, y) \le C_1 t^{-n/2} \exp\left(-C_2 d(y, \partial D)^2/t\right)$$
 (2)

for  $0 < t \le \min(T, 2C_2 d(y, \partial D)^2/n)$ .

*Proof.* The first inequality (1) is due to [Do] (see also [BS]). The second inequality is a consequence of the maximum principle (see [C] and [D]).

PROPOSITION 2.1. If the family D satisfies the property (P) then

$$\lim_{j \to \infty} \text{vol}(D_j)^{-1} \int_{D_j} (k(t, x, x) - k_{D_j}(t, x, x)) dx = 0.$$

*Proof.* Let t > 0, and take constants  $C_1$  and  $C_2$  in (1) for T = t. We have

$$vol (D_{j})^{-1} \int_{D_{j}} (k(t, x, x) - k_{D_{j}}(t, x, x)) dx$$

$$= vol (D_{j})^{-1} \int_{\partial_{h}D_{j}} (k(t, x, x) - k_{D_{j}}(t, x, x)) dx$$

$$+ vol (D_{j})^{-1} \int_{D_{j} \setminus \partial_{h}D_{j}} (k(t, x, x) - k_{D_{j}}(t, x, x)) dx.$$

In view of Lemma 2.2, (1), the first term is estimated from above by

$$C_1 t^{-n/2} \operatorname{vol}(\partial_h D_i) / \operatorname{vol}(D_i),$$

which tends to zero as  $j \uparrow \infty$ . Take h with  $t \le 2C_2h^2/n$ . Then, for  $x \in D_j \setminus \partial_h D_j$ , one has  $t \le 2C_2 d(x, \partial D_j)^2/n$ , so that, by Lemma 2.2, (2) the second term is estimated from above by

$$C_1 t^{-n/2} \exp(-C_2 h^2/t)$$
.

By letting h go to infinity, we get the assertion.

**PROPOSITION** 2.2. If  $\mathcal{D}$  satisfies the property (P), then one has, for every  $\Gamma$ -periodic continuous function f, that

$$\lim_{j\to\infty}\operatorname{vol}(D_j)^{-1}\int_{D_j}f(x)\,dx=\operatorname{vol}(\mathscr{F})^{-1}\int_{\mathscr{F}}f(x)\,dx.$$

*Proof.* Put  $E_j = \{ \sigma \in \Gamma; (D_j \setminus \partial_h D_j) \cap \sigma \overline{\mathscr{F}} \neq \emptyset \}$ , and

$$D'_j = \bigcup_{\sigma \in E_j} \sigma \bar{\mathscr{F}}.$$

It is clear that  $(D_j \setminus \partial_h D_j) \subset D_j'$ . We show that, if  $h \ge \text{diam}(\mathscr{F})$ , then  $D_j' \subset D_j$ . Let  $x \in (D_j \setminus \partial_h D_j) \cap \sigma \mathscr{F}$ . since  $d(x, \partial D_j) \ge h$ , we find  $\partial D_j \cap B_h(x) = \emptyset$ , where  $B_h(x) = \{z \in X; d(x, z) \le h\}$ . From the connectedness of  $B_h(x)$ , it follows that  $B_h(x) \subset D_j$ . Since  $h \ge \text{diam}(\mathscr{F})$ , we have  $\sigma \mathscr{F} \subset B_h(x) \subset D_j$ .

We now find

$$\operatorname{vol}(D_{j})^{-1} \int_{D_{j}} f(x) \, dx = \operatorname{vol}(D_{j})^{-1} \int_{D_{j}} f(x) \, dx + \operatorname{vol}(D_{j})^{-1} \int_{D_{j} \setminus D_{j}'} f(x) \, dx$$

$$= \frac{\operatorname{vol}(D_{j}')}{\operatorname{vol}(D_{j})} \frac{1}{\operatorname{vol}(D_{j}')} \int_{D_{j}} f(x) \, dx$$

$$+ \frac{\operatorname{vol}(D_{j} \setminus D_{j}')}{\operatorname{vol}(D_{j})} \frac{1}{\operatorname{vol}(D_{j} \setminus D_{j}')} \int_{D_{j} \setminus D_{j}} f(x) \, dx.$$

Since  $D_j \setminus D_j' \subset \partial_h D_j$ , we have  $\lim_{j \to \infty} \text{vol}(D_j \setminus D_j')/\text{vol}(D_j) = 0$  and  $\lim_{j \to \infty} \text{vol}(D_j')/\text{vol}(D_j) = 1$ . In view of the  $\Gamma$ -periodicity of f, we find

$$\lim_{j \to \infty} \operatorname{vol}(D_j)^{-1} \int_{D_j} f(x) \, d(x) = \lim_{j \to \infty} \operatorname{vol}(D'_j)^{-1} \int_{D'_j} f(x) \, dx$$
$$= \operatorname{vol}(\mathscr{F})^{-1} \int_{\mathscr{F}} f(x) \, dx.$$

We shall make use of the following genral lemma to complete the proof of Theorem 1.1.

LEMMA 2.3 (cf. [Sh]). Let  $\{\varphi_j(\lambda)\}_{j=1}^{\infty}$  be a sequence of non-decreasing functions with  $\varphi_j(\lambda) = 0$  for  $\lambda \leq c$ , where c is a constant not depending on j. Suppose that there exists a function C(t), not depending on j such that

$$\Phi_j(t) := \int e^{-\lambda t} d\varphi_j(t) \le C(t),$$

and

$$\lim_{j\to\infty}\Phi_j(t)=\int e^{-\lambda t}\,d\varphi(\lambda),$$

where  $\varphi$  is a non-decreasing function. Then  $\lim_{j\to\infty} \varphi_j(\lambda) = \varphi(\lambda)$  at all points of continuity of  $\varphi(\lambda)$ .

We apply this lemma to

$$\varphi_j(\lambda) = \text{vol } (D_j)^{-1} \varphi_{D_j}(\lambda),$$
  
$$\varphi(\lambda) = \text{vol } (\Gamma \backslash X)^{-1} \Phi_{\Gamma}(\lambda).$$

Since the first eigenvalue of  $H_{D_j}$  is not less than min q(x), we observe that  $\varphi_j(\lambda) = 0$  for  $\lambda < \min q(x)$ . We also find that

$$\Phi_{j}(\lambda) = \operatorname{vol}(D_{j})^{-1} \int_{D_{j}} k_{D_{j}}(t, x, x) dx$$

$$\leq \operatorname{vol}(D_{j})^{-1} \int_{D_{j}} k(t, x, x) dx$$

$$\leq \sup_{x \in X} k(t, x, x) =: C(t),$$

where we should note that the function k(t, x, x) is  $\Gamma$ -periodic with respect to the variable x. By Proposition 2.1,

$$\lim_{j \to \infty} \Phi_j(t) = \lim_{j \to \infty} \operatorname{vol}(D_j)^{-1} \int_{D_j} k_{D_j}(t, x, x) dx$$
$$= \lim_{j \to \infty} \operatorname{vol}(D_j)^{-1} \int_{D_j} k(t, x, x) dx.$$

Using again  $\Gamma$ -periodicity of k(t, x, x), together with Proposition 2.2, we have

$$\lim_{j \to \infty} \Phi_j(t) = \text{vol}(\mathscr{F})^{-1} \int_{\mathscr{F}} k(t, x, x) dx$$

$$= \text{vol}(\mathscr{F})^{-1} \int e^{-\lambda t} d\Phi_{\Gamma}(\lambda)$$

$$= \int e^{-\lambda t} d\varphi(\lambda)$$

as desired.

## §3. Manifolds with amenable group actions

In this section, we shall prove Proposition 1.1 in a slightly strong form. For this, we recall the Følner's characterization of amenability. Let  $\Gamma$  be a finitely generated group with a fixed finite set A of generators.

PROPOSITION 3.1 (Følner [F] and [Ad]).  $\Gamma$  is amenable if and only if, for every positive  $\varepsilon$ , there exists a non-empty finite set E such that

$$|EA\setminus E|\leq \varepsilon |E|,$$

where |E| denotes the cardinality of the set E.

We first assume that a manifold X with compact quotient  $\Gamma \setminus X$  has a family  $\{D_i\}$  satisfying the property (P). Fixing a fundamental domain  $\mathcal{F}$ , we put

$$A = \{ a \in \Gamma \colon a\overline{\mathcal{F}} \cap \overline{\mathcal{F}} \neq \emptyset \}.$$

The finite set A generates  $\Gamma$ . Taking a number  $h > 2 \cdot \text{diam } (\mathcal{F})$ , we set

$$E_i = \{ \gamma \in \Gamma; \gamma \bar{\mathscr{F}} \cap (D_i \setminus \partial_h D_i) \neq \emptyset \}.$$

Let  $\sigma = \gamma \cdot a \in E_j A$   $(\gamma \in E_j, a \in A)$ . We shall prove that  $\sigma \overline{\mathscr{F}} \subset D_j$ . For this, let  $z \in \gamma \overline{\mathscr{F}} \cap (D_j \setminus \partial_h D_j)$ . We then have  $B_h(z) \subset D_j$  as before. Since

$$\sigma \overline{\mathscr{F}} \cap \gamma \overline{\mathscr{F}} = \gamma (a \overline{\mathscr{F}} \cap \overline{\mathscr{F}}) \neq \emptyset,$$

there exists an element  $y \in \sigma \overline{\mathscr{F}} \cap \gamma \overline{\mathscr{F}}$ , and hence, for every  $x \in \sigma \overline{\mathscr{F}}$ , one has

$$d(x, z) \le d(x, y) + d(y, z) \le 2 \cdot \operatorname{diam}(\mathscr{F}) < h,$$

which implies that  $\sigma \overline{\mathscr{F}} \subset B_h(z)$ , and hence  $\sigma \overline{\mathscr{F}} \subset D_j$ . We now observe

$$|E_{j}A \setminus E_{j}|/|E_{j}| = \frac{|E_{j}A| \operatorname{vol}(\mathscr{F})}{|E_{j}| \operatorname{vol}(\mathscr{F})} - 1 \leq \frac{\operatorname{vol}(D_{j})}{\operatorname{vol}(D_{j} \setminus \partial_{h}D_{j})} - 1$$
$$= \frac{\operatorname{vol}(\partial_{h}D_{j})}{\operatorname{vol}(D_{j})} \left(1 - \frac{\operatorname{vol}(\partial_{h}D_{j})}{\operatorname{vol}(D_{j})}\right)^{-1},$$

which goes to zero as  $j \uparrow \infty$ . Hence  $\Gamma$  is amenable by Følner's criterion.

Next we suppose that  $\Gamma$  is amenable. Using a smooth triangulation of the orbifold  $\Gamma \setminus X$ , we may lift up *n*-simplices one by one to X to obtain a connected polyhedral fundamental domain  $\mathscr{F}$ . The finite set  $A = \{\sigma \in \Gamma; \sigma \overline{\mathscr{F}} \cap \overline{\mathscr{F}} \neq \emptyset\}$  is symmetric and contains the unit element. We associate the Cayley graph  $\mathscr{C}(\Gamma, A)$ ; the set of vertices being  $\Gamma$  and the set of edges being  $\{(\gamma, \sigma) \in \Gamma \times \Gamma; \gamma^{-1}\sigma \in A\}$ . We denote by  $d_A$  the distance function on  $\Gamma$  associated with the graph  $\mathscr{C}(\Gamma, A)$ . A subset E in  $\Gamma$  will be called connected if, for any two vertices in E, there exists a path in  $\mathscr{C}(\Gamma, A)$  joining those vertices and consisting of vertices in E. By use of Theorem 4 in [Ad], there is a family  $\{E_i\}_{i=1}^{\infty}$  of connected subsets of  $\Gamma$  such that

$$\bigcup_{j=1}^{\infty} E_j = \Gamma, \qquad E_j \subset E_j \cdot A \subset E_{j+1} \quad \text{and}$$
$$|E_i \cdot A^j \setminus E_i| \le |E_i|/j|A|^j \quad \text{for every } j.$$

We put  $F_j = \bigcup_{\gamma \in E_j} \gamma \mathcal{F}$  and  $F'_j = \bigcup_{\gamma \in E_j \cdot A} \gamma \bar{F}$ , which are connected by the choice of A and the connectedness of  $E_j$ . It should be noted that there exists a positive  $\varepsilon$  such that the  $\varepsilon$ -neighborhood of  $F_j$  is contained in  $F'_j$ . Thus we may make a uniform

regularization  $D_j$  of  $F_j$  satisfying  $F_j \subset D_j \subset \bar{D}_j \subset F_j'$  (see [B]). It is clear that  $\bigcup_{j=1}^{\infty} D_j = X$  and  $\bar{D}_j \subset D_{j+1}$ . Our goal is to show that  $\{D_j\}_{j=1}^{\infty}$  satisfies the property (P). Let  $x_0 \in \mathcal{F}$ . Since the map  $f: \Gamma \to X$ ,  $f(\gamma) = \gamma x_0$ , is a rough isometry (Kanai [K]), we have

$$d_A(\gamma, \mu) \le c_1 d(\gamma x_0, \mu x_0) + c_2$$

with suitable constants  $c_1 > 0$  and  $c_2 \ge 0$ .

LEMMA 3.1. If  $h \le (j - c_2)/c_1 - 2 \cdot \text{diam}(\mathcal{F})$ , then the thick boundary  $\partial_h D_j$  is contained in the set

$$\partial^j F_i' = \bigcup \{ \mu \sigma \mathscr{F} ; \sigma \in A, \mu \in E_i, \text{ and there is } \gamma \in A^j \text{ with } \mu \gamma \notin E_i \}.$$

*Proof.* Suppose x is contained in  $\partial_h D_j \cap \mu \overline{\mathscr{F}}$  for some  $\mu \in E_j$ . Since  $F_j \subset D_j$  there is  $y \in \overline{X \setminus F_j}$  with  $d(x, y) \leq h$ . Choose  $\rho \notin E_j$  so that  $y \in \rho \overline{\mathscr{F}}$ . Then  $d(\mu x_0, \rho x_0) \leq h + 2 \cdot \text{diam } (\mathscr{F})$ , hence  $d_A(\mu, \rho) \leq j$  and  $\partial_h D_j \cap F_j \subset \partial^j F_j$ , where

$$\partial^j F_i = \bigcup \{ \mu \mathscr{F} \mid \mu \in E_i \text{ and there is } \gamma \in A^j \text{ with } \mu \gamma \notin E_i \}.$$

If  $\gamma \in E_i A \setminus E_i$ , it is clear that  $\gamma \mathscr{F} \subset \partial^j F_i'$  (since  $A \subset A^j$ ), therefore

$$\partial_h D_i \subset (\partial_h D_i \cap F_i) \cup (F'_i \setminus F_i) \subset \partial^j F'_i$$
.

We now show that the family  $\{D_j\}_{j=1}^{\infty}$  satisfies the property (P). By the definition of  $\partial^j F_j$  and  $\partial^j F_j'$  we have

$$\operatorname{vol}(\partial^{j} F_{j}') \leq |A| \cdot \operatorname{vol}(\partial^{j} F_{j})$$

$$= |A| \cdot \operatorname{vol}(\mathscr{F}) \cdot |\{\mu \in E_{j} \mid \text{there is } \gamma \in A^{j} \text{ with } \mu \gamma \notin E_{j}\}|$$

$$\leq |A| \cdot \operatorname{vol}(\mathscr{F}) \sum_{\gamma \in A^{j}} |E_{j} \setminus E_{j} \gamma^{-1}|$$

$$= |A| \cdot \operatorname{vol}(\mathscr{F}) \sum_{\gamma \in A^{j}} |E_{j} \gamma \setminus E_{j}|$$

$$\leq |A| \cdot \operatorname{vol}(\mathscr{F}) \cdot |A^{j}| \cdot |E_{j} A^{j} \setminus E_{j}|$$

$$\leq \operatorname{vol}(\mathscr{F}) \cdot |E_{j}|/j$$

$$= \operatorname{vol}(F_{j})/j.$$

Therefore we get, for every h > 0, that

$$\operatorname{vol}(\partial_h D_i)/\operatorname{vol}(D_i) \leq \operatorname{vol}(\partial^j F_i)/\operatorname{vol}(F_i) \leq 1/j \to 0.$$

Summarizing up, we obtain

**PROPOSITION** 3.2. If  $\Gamma$  is amenable, then there exists an expanding family  $\mathcal{D} = \{D_j\}$  of bounded open domains with smooth boundaries satisfying the following conditions:

- (1)  $\mathcal{D}$  has the property (P),
- (2) the boundary  $\partial D_j$  has a uniformly bounded second fundamental form  $h_j$ . More precisely, there exists positive constant c not depending on j with  $-cg \le h_j \le cg$ , where g denotes the Riemannian metric on X.

A group of subexponential growth is amenable (see [B]). In this case, we may construct a family  $\mathcal{D} = \{D_j\}$  satisfying the conditions in the above proposition by using the following property on concentric geodesic balls.

LEMMA 3.2. Suppose that  $\Gamma$  is of subexponential growth. For an arbitrary point x in X, there is a sequences of positive numbers  $\{R_i\}_{i=1}^{\infty}$  such that

- (1)  $R_i \uparrow \infty$ ,
- (2)  $\lim_{i \to \infty} \text{vol}(B_{R_i}(x))/\text{vol}(B_{R_i h}(x)) = 1$  for every h > 0.

(cf. [Ad]).

# §4. Hyperoblic spaces

We now consider the density of states associated with the Laplacian on the hyperbolic space  $X = \mathbb{H}^n$ . The manifold  $\mathbb{H}^n$  is a typical example of a manifold with a non-amenable discontinuous transformation group.

THEOREM 4.1. Let  $\mathcal{D} = \{D_j\}$  be a family of concentric geodesic balls in  $\mathbb{H}^n$ . Then one has

$$\operatorname{vol}(\mathscr{F})^{-1} \int e^{-\lambda t} d\Phi_{\Gamma}(\lambda) > \lim_{j \to \infty} \sup \operatorname{vol}(D_j)^{-1} \int e^{-\lambda t} d\varphi_{D_j}(\lambda).$$

In particular, vol  $(\Gamma \setminus X)^{-1}\Phi_{\Gamma} \neq \varphi_{\mathscr{D}}$ .

*Proof.* Since  $\mathbb{H}^n$  is a homogeneous Riemannian manifold, k(t, x, x) does not depend on the variable x, so that we write

$$k(t) = k(t, x, x).$$

We then find

$$\operatorname{vol}(\mathscr{F})^{-1} \int e^{-\lambda t} d\Phi_{\Gamma}(\lambda) - \operatorname{vol}(D_{j})^{-1} \int_{D_{j}} k_{D_{j}}(t, x, x) dx$$

$$= k(t) - \operatorname{vol}(D_{j})^{-1} \int_{D_{j}} k_{D_{j}}(t, x, x) dx$$

$$= \operatorname{vol}(D_{j})^{-1} \int_{D_{j}} (k(t) - k_{D_{j}}(t, x, x)) dx$$

$$\geq \operatorname{vol}(D_{j})^{-1} \int_{\partial_{D_{j}}} (k(t) - k_{D_{j}}(t, x, x)) dx,$$

where we have used the fact that  $k_D(t, x, y) \le k(t)$ .

To complete the proof, we need the following lemma.

LEMMA 4.1. For a fixed t > 0, there exists a positive h such that

$$k_D(t, x, x) \le k(t)/2$$

for every geodesic ball D and every  $x \in \partial_h D$ .

*Proof.* Choose a unit speed geodesic  $C: \mathbb{R} \to X$ , and consider the horoball  $H = \bigcup_{\tau > 0} B_{\tau}(c(\tau))$ . Let  $k_H(t, x, y)$  denote the Dirichlet heat kernel function for the horoball. Since  $\lim_{x \to \partial H} k_H(t, x, x) = 0$ , it follows that there exists a positive h such that, for a positive  $\delta$  with dist  $(c(\delta), \partial H) = \delta \leq h$ .

$$k_H(t, c(\delta), c(\delta)) \le k(t)/2.$$

Let  $x \in \partial_h D$ . Since one can find an isometry f on  $\mathbb{H}^n$  such that  $f(D) = B_{\tau}(c(\tau))$ ,  $\tau > 0$ , and  $f(x) = c(\delta)$  for some  $\delta \le h$ . Hence we have, by the domain monotonicity of the Dirichlet heat kernel,

$$k_D(t, x, x) = k_{B_{\tau}(c(\tau))}(t, c(\delta), c(\delta))$$

$$\leq k_H(t, c(\delta), c(\delta)) \leq k(t)/2,$$

as desired.

Applying the above lemma, we get

$$\operatorname{vol}(\mathscr{F})^{-1} \int e^{-\lambda t} d\Phi_{\Gamma}(\lambda) - \operatorname{vol}(D_{j})^{-1} \int_{D_{j}} k_{D_{j}}(t, x, x) dx$$

$$\geq \frac{k(t)}{2} \operatorname{vol}(\partial_{h} D_{j}) / \operatorname{vol}(D_{j}).$$

If  $r_j$  denotes the radius of  $D_j$ , one has vol  $(D_j) = e^{(n-1)r_j}$ , so that the last term is written as

$$\frac{k(t)}{2}(1-e^{-(n-1)h})>0.$$

This completes the proof of Theorem 4.1.

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