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# On the Covariant Differential of an Almost Hermitian Structure

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This paper deals with the covariant differential  $\nabla J$  of an almost hermitian structure  $J$  on a Riemannian manifold  $M$  with metric  $g$ . The connection with respect to which  $\nabla J$  is defined is the Riemannian connection on  $M$ . The extension of the notion of antisymmetrization and symmetrization from tensor fields of type  $(0, q)$  to those of type  $(p, q)$  for  $p > 0$  allows us to decompose  $\nabla J$  in an antisymmetric part  $A$  and a symmetric part  $S$ . In this way, we find new formulas which have the feature of stressing the relationship between  $\nabla J$  and the torsion  $\tau$  of  $J$  as well as the fundamental 2-form  $\omega$  or better, its exterior derivative  $d\omega$ . The results are essentially based first on the Palais formula which gives the exterior derivative of a  $q$ -form via Lie product and covariant derivative and second on a theorem ([2], p. 149) which gives the exterior derivative  $d\alpha$  of a  $q$ -form  $\alpha$  as the antisymmetric part of the covariant differential  $\nabla\alpha$  of  $\alpha$ , provided the connection has vanishing torsion. As an application of our formulas we give a characterization of so called nearly Kähler manifolds ([1]) via the fundamental 2-form  $\omega$ . We also give a very simple proof of the characterization of a Kähler manifold given by the vanishing of  $\nabla J$  or of  $d\omega$  and  $\tau$ . We finally prove a lemma which gives a nice interpretation of the torsion of  $J$  when the fundamental 2-form is closed, that is, in the case of an almost Kähler manifold.

## §1. The Covariant Differential of a Tensor Field

Let  $t$  be a given tensor field of type  $(p, q)$  on a  $C^\infty$  manifold  $M$ . We shall simply write  $t \in T_M(p, q)$  or  $t \in T(p, q)$ .

Suppose there is also a linear connection  $\nabla$  given on  $M$ . Then, as in [2], we can define the covariant differential of  $t$  as the tensor field  $\nabla t \in T(p, q+1)$  defined by

$$\nabla t(X_1, \dots, X_q, X) = (\nabla_X t)(X_1, \dots, X_q) \quad (1.1)$$

where  $\nabla_X t$  denotes the covariant derivative of the tensor field  $t$  and  $X_1, \dots, X_q, X$  are in the Lie algebra  $\mathfrak{X}(M)$  of vector fields on  $M$ .

Because  $\nabla_X$  is a derivation commuting with every contraction, we have

**THEOREM 1.1** ([2], p. 124). *If  $t \in T(p, q)$  then, for  $X_i, X \in \mathfrak{X}(M)$*

$$\nabla t(X_1, \dots, X_q, X) = \nabla_X(t(X_1, \dots, X_q)) - \sum_{k=1}^q t(X_1, \dots, \nabla_{X_i} X_k, \dots, X_q).$$

EXAMPLE 1. Take  $t=g$  a Riemann metric on  $M$ .  $g$  is in  $T(0, 2)$ , so  $\nabla g$  is in  $T(0, 3)$  and for  $X, Y, Z, \in \mathfrak{X}(M)$  we have

$$\nabla g(X, Y, Z) = (\nabla_Z g)(X, Y) = Zg(X, Y) - g(\nabla_Z X, Y) - g(X, \nabla_Z Y).$$

One of the features of a Riemannian metric  $g$  is to be parallel i.e.  $\nabla g = 0$ . For convenience we shall write  $\langle X, Y \rangle$  instead of  $g(X, Y)$ .

EXAMPLE 2. Take  $t=J$  an almost complex structure on  $M$ . This is a tensor field of type  $(1, 1)$  whose square  $J^2$  equals minus the identity.  $\nabla J$  is in  $T(1, 2)$  and

$$\nabla J(X, Y) = (\nabla_Y J)X = \nabla_Y(JX) - J\nabla_Y X.$$

One of the features of a Kähler structure  $J$  on  $M$  is to be parallel with respect to  $\nabla$ :  $\nabla J = 0$ .

## §2. Extension of Antisymmetrization and Symmetrization

It is known [1] how to define for a covariant tensor field  $t$  in  $T(0, q)$  the *alternation*  $At$  of  $t$ . It is a tensor field of the same type defined by

$$(At)(X_1, \dots, X_q) = \frac{1}{q!} \sum_{\pi \in P_q} \varepsilon(\pi) t(X_{\pi(1)}, \dots, X_{\pi(q)}) \quad (2.1)$$

where  $P_q$  is the group of permutations of  $\{1, 2, \dots, q\}$  and  $\varepsilon(\pi)$  is the signe of the permutation  $\pi$ . Notice that  $At$  is a skew-symmetric tensor field and  $t$  is skew-symmetric if and only if  $At = t$ .

On the other hand one also defines for  $t \in T(0, q)$  the *symmetrization*  $St$  of  $t$  by

$$St(X_1, \dots, X_q) = \frac{1}{q!} \sum_{\pi \in P_q} t(X_{\pi(1)}, \dots, X_{\pi(q)}). \quad (2.2)$$

Here is  $St$  a symmetric tensor field and  $t$  is symmetric if and only if  $St = t$ .

We now make the straightforward extension of the above notions to tensor fields of type  $(p, q)$  for  $p > 0$ . In this case  $t(X_1, \dots, X_q)$  is no longer a real function on  $M$ , but a  $p$ -contravariant tensor field. Nevertheless, all algebraic operations needed for definitions (2.1) and (2.2) still make sense in the module of  $p$ -contravariant tensor fields.

EXAMPLE 1. Take  $t=\tau \in T(1, 2)$  the torsion of an almost complex structure  $J$  defined by

$$\tau(X, Y) = [X, Y] + J[JX, Y] + J[X, JY] - [JX, JY].$$

Here we have  $A\tau = \tau$  and  $S\tau = 0$ .

EXAMPLE 2. Take  $t = K \in T(1, 3)$  where  $K(X, Y, Z) = R(X, Y)Z$  and  $R \in T(1, 3)$  is defined by  $R(X, Y) = [\nabla_X, \nabla_Y] - \nabla_{[X, Y]}$ .

Because  $R(X, Y) = -R(Y, X)$  we have

$$(AK)(X, Y, Z) = \frac{1}{3}\mathfrak{s}(K(X, Y, Z))$$

( $\mathfrak{s}$  means cyclic sum) and  $SK = 0$ . Furthermore if the torsion of  $\nabla$  vanishes, we have  $AK = 0$  in view of the first Bianchi identity.

### §3. Almost Hermitian Manifolds

From now on, we assume that  $(M, J, g)$  is an almost hermitian manifold, that is, the almost complex structure  $J$  on  $M$  gives in each point  $p$  of  $M$  an isometry of  $T_p(M)$ , the tangent space at  $M$  in  $p$ , i.e.

$$\langle JX, JY \rangle = \langle X, Y \rangle \quad \forall X, Y \in T_p(M). \quad (3.1)$$

Denoting by  $\nabla$  the Riemannian connection on  $M$ , we can compute  $\nabla J$  which is a tensor field of type  $(1, 2)$  and we can write the following decomposition

$$\nabla J = A(VJ) + S(VJ). \quad (3.2)$$

Needless to say, this is only possible because we are in  $T(p, q)$  with  $q = 2$ . Example 2 above shows what can happen with  $q \neq 2$ .

For convenience we shall write  $A$  for  $A(\nabla J)$  and  $S$  for  $S(\nabla J)$  and establish several formulas relating  $A$  and/or  $S$  with various tensor fields one can define on an almost hermitian manifold.

#### I. The Torsion $\tau$ and $A$

THEOREM 3.1. *If  $A$  denotes the antisymmetric part of the covariant differential  $\nabla J$  of the almost hermitian structure  $J$  on  $M$ , then*

$$\frac{1}{2}J\tau(X, Y) = A(X, Y) - A(JY, JY). \quad (3.3)$$

*Proof.* By definition  $A(X, Y) = \frac{1}{2}\{\nabla J(X, Y) - \nabla J(Y, X)\}$ . So

$$2\{A(X, Y) - A(JX, JY)\} = \nabla J(X, Y) - \nabla J(Y, X) - \nabla J(JX, JY) + \nabla J(JY, JX).$$

In view of (1.1), the right hand side becomes

$$(\nabla_Y J)X - (\nabla_X J)Y - (\nabla_{JY} J)(JX) + (\nabla_{JX} J)(JY).$$

which is, according to Theorem 1.1, the same as

$$\nabla_Y JX - J\nabla_Y X - \nabla_X JY + J\nabla_X Y + \nabla_{JY} X + J\nabla_{JY} JX - \nabla_{JX} Y - J\nabla_{JX} JY. \quad (3.4)$$

The Riemannian connection having torsion zero, we have  $\nabla_X Y - \nabla_Y X = [X, Y]$  and multiplication of (3.4) by  $J$  gives (3.3).

From Theorem 3.1 we trivially get the following

**COROLLARY 3.2.** *If the covariant differential  $\nabla J$  of an almost hermitian structure  $J$  is symmetric then  $J$  is integrable.*

## II. The Fundamental 2-form $\omega$ and $A$

The fundamental 2-form  $\omega$  on an almost hermitian manifold  $(M, J, g)$  is defined by

$$\omega(X, Y) = \langle JX, Y \rangle \quad \text{for } X, Y \in \mathfrak{X}(M). \quad (3.5)$$

The torsion of the Riemann connection being zero, we know ([2], chap. III) that  $d\omega = A(\nabla\omega)$ . For vector fields  $X, Y$ , and  $Z$  on  $M$ , we have therefore:

$$6d\omega(X, Y, Z) = 6A(\nabla\omega)(X, Y, Z) = \mathfrak{s}(\nabla\omega(X, Y, Z)) - \mathfrak{s}(\nabla\omega(Y, X, Z)) \quad (3.6)$$

where  $\mathfrak{s}$  denotes cyclic sum. By definition of the covariant differential and (3.5) we have

$$\begin{aligned} \nabla\omega(X, Y, Z) &= (\nabla_Z \omega)(X, Y) = Z\omega(X, Y) - \omega(\nabla_Z X, Y) - \omega(X, \nabla_Z Y) \\ &= \langle \nabla_Z JX, Y \rangle + \langle JX, \nabla_Z Y \rangle - \langle J\nabla_Z X, Y \rangle - \langle JX, \nabla_Z Y \rangle \\ &= \langle \nabla_Z JX, Y \rangle - \langle J\nabla_Z X, Y \rangle = \langle (\nabla_Z J)X, Y \rangle = \langle \nabla J(X, Z), Y \rangle. \end{aligned}$$

Hence

$$\nabla\omega(X, Y, Z) = \langle \nabla J(X, Z), Y \rangle \quad (3.7)$$

In view of (2.1) and (3.2) we have the following

**THEOREM 3.3.** *On an almost hermitian manifold  $(M, J, g)$  the exterior differential  $d\omega$  of the fundamental 2-form  $\omega$  and the antisymmetric part of the covariant differential  $\nabla J$  are related by the formula*

$$3d\omega(X, Y, Z) = -\mathfrak{s}(\langle A(X, Y), Z \rangle). \quad (3.8)$$

Together with Corollary 3.2 this result implies

**COROLLARY 3.4.** *If the covariant differential  $\nabla J$  of an almost hermitian structure*

*J is symmetric, then the fundamental 2-form  $\omega$  is closed and J is integrable, i.e.  $(M, J, g)$  is a Kähler manifold.*

There is another interesting relation between  $d\omega$ ,  $A$  and  $\nabla J$ ; namely, we have

**THEOREM 3.5.** *On any almost hermitian manifold  $(M, J, q)$  we have for the fundamental 2-form  $\omega$ :*

$$3d\omega(X, Y, Z) = -2\langle A(X, Y), Z \rangle + \langle \nabla J(X, Z), Y \rangle. \quad (3.9)$$

*Proof.* By the Palais formula, one has for any 2-form  $\omega$

$$\begin{aligned} 3d\omega(X, Y, Z) &= X\omega(Y, Z) - Y\omega(X, Z) + Z\omega(X, Y) \\ &\quad - \omega([X, Y], Z) + \omega([X, Z], Y) - \omega([Y, Z], X). \end{aligned}$$

By the definition of  $\omega$  and because  $g$  is parallel with respect to  $\nabla$ , together with the fact, once again, that the torsion of  $\nabla$  vanishes, we get the stated result.

### III. A Relation Between $A$ and $S$ and an Identity for $S$

From Theorem 3.3 and Theorem 3.5 above we deduce the

**COROLLARY 3.6.** *If  $A$  (resp.  $S$ ) is the antisymmetric (resp. symmetric) part of the covariant differential  $\nabla J$  of an almost hermitian structure then, for vector fields  $X, Y, Z$  on  $M$ , we have*

$$\langle S(X, Y), Z \rangle = \langle A(X, Z), Y \rangle + \langle A(Y, Z), X \rangle. \quad (3.10)$$

*Proof.* From (3.2), (3.8) and (3.9) one has

$$\mathfrak{s}(\langle A(X, Y), Z \rangle) = 2\langle A(X, Y), Z \rangle + \langle A(X, Z), Y \rangle + \langle S(X, Z), Y \rangle.$$

But  $A$  is antisymmetric, so  $A(X, Z) = -A(Z, X)$  and (3.10) follows by permuting  $Y$  and  $Z$ .

The antisymmetry of  $A$  has another consequence:

**COROLLARY 3.7.** *The symmetric part  $S$  of the covariant differential  $\nabla J$  of an almost hermitian structure satisfies the following identity*

$$\mathfrak{s}(\langle S(X, Y), Z \rangle) = 0 \quad (3.11)$$

where  $\mathfrak{s}$  denotes cyclic sum.

*Proof.* Write the left hand side of (3.11) with (3.10) and use the antisymmetry of  $A$ .

#### §4. A few Remarks

1. Because of (3.11) we can rewrite Theorem 3.3 in the following form:

**THEOREM 3.3'.** *On an almost hermitian manifold  $(M, J, g)$  the exterior differential  $d\omega$  of the fundamental 2-form  $\omega$  and the covariant differential  $\nabla J$  of  $J$  are related by the formula*

$$3d\omega(X, Y, Z) = -s(\langle \nabla J(X, Y), Z \rangle). \quad (4.1)$$

2. An almost hermitian manifold  $(M, J, g)$  for which  $S=0$  is already known [1] as *nearly Kähler manifold*. An alternative condition is given by the following

**THEOREM 4.1.** *A nearly Kähler manifold  $(M, J, g)$  is characterized by the condition*

$$\nabla\omega = d\omega. \quad (4.2)$$

*Proof.* We have to show that this condition is equivalent to  $S=0$ . Suppose (4.2) is true. Then  $\nabla\omega$  is antisymmetric. By (3.7)  $\nabla\omega(X, Y, Z) = \langle \nabla J(X, Z), Y \rangle$  and  $\nabla\omega(Z, Y, X) = \langle \nabla J(Z, X), Y \rangle = -\langle \nabla J(X, Z), Y \rangle$  which implies that  $\nabla J$  itself is antisymmetric, i.e.  $S=0$ .

On the other hand,  $S=0$  implies  $\nabla J = A$  and by (3.10)  $\langle \nabla J(X, Z), Y \rangle + \langle \nabla J(Y, Z), X \rangle = 0$  which in turn gives by antisymmetry of  $\nabla J$ :  $\langle \nabla J(Z, X), Y \rangle + \langle \nabla J(Z, X), Y \rangle = 0$  or with (3.7),  $\nabla\omega(Z, Y, X) + \nabla\omega(Z, X, Y) = 0$ . But (even if  $S \neq 0$ ) one has  $\nabla\omega(X, Y, Z) + \nabla\omega(Y, X, Z) = 0$ , and (3.6) gives the result.

3. Of the antisymmetric part  $A$  and the symmetric part  $S$  of  $\nabla J$ , the former plays the most important role. It allows us to give as an application a very simple proof of the following theorem. Compare with [3] (chap. IX).

**THEOREM 4.2.** *An almost hermitian manifold  $(M, J, g)$  is a Kähler manifold (i.e.  $\tau=0$  and  $d\omega=0$ ) if and only if the covariant differential  $\nabla J$  vanishes.*

*Proof.* If  $\nabla J=0$  then  $A=0$  which implies  $\tau=0$  by (3.3) and  $d\omega=0$  by (3.8). On the other hand, from Lemma 4.3 below and  $d\omega=0$  we get  $A=0$  and by (3.10),  $S=0$ .

**LEMMA 4.3.** *For an almost Kähler manifold, the antisymmetric part  $A$  of the covariant differential  $\nabla J$  is essentially the torsion of  $J$ : more precisely we have*

$$4A = J\tau.$$

*Proof.*  $d\omega=0$  implies by (3.9)

$$(\nabla J(X, Z), Y) = 2\langle A(X, Y), Z \rangle.$$

Substituting  $JX$  to  $X$  and  $JY$  to  $Y$  and adding, we get

$$\langle \nabla J(X, Z), Y \rangle + \langle \nabla J(JX, Z), JY \rangle = 2 \langle A(X, Y) + A(JX, JY), Z \rangle. \quad (4.3)$$

But the left hand side vanishes because  $\nabla J(JX, Y) = -J\nabla J(X, Y)$  as it is easy to see and  $J$  is an isometry. To get the desired result one just has to add (4.3) to (3.3).

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